

Quantum devices for noise- induced switching, signal detection and energy harvesting

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Systems*

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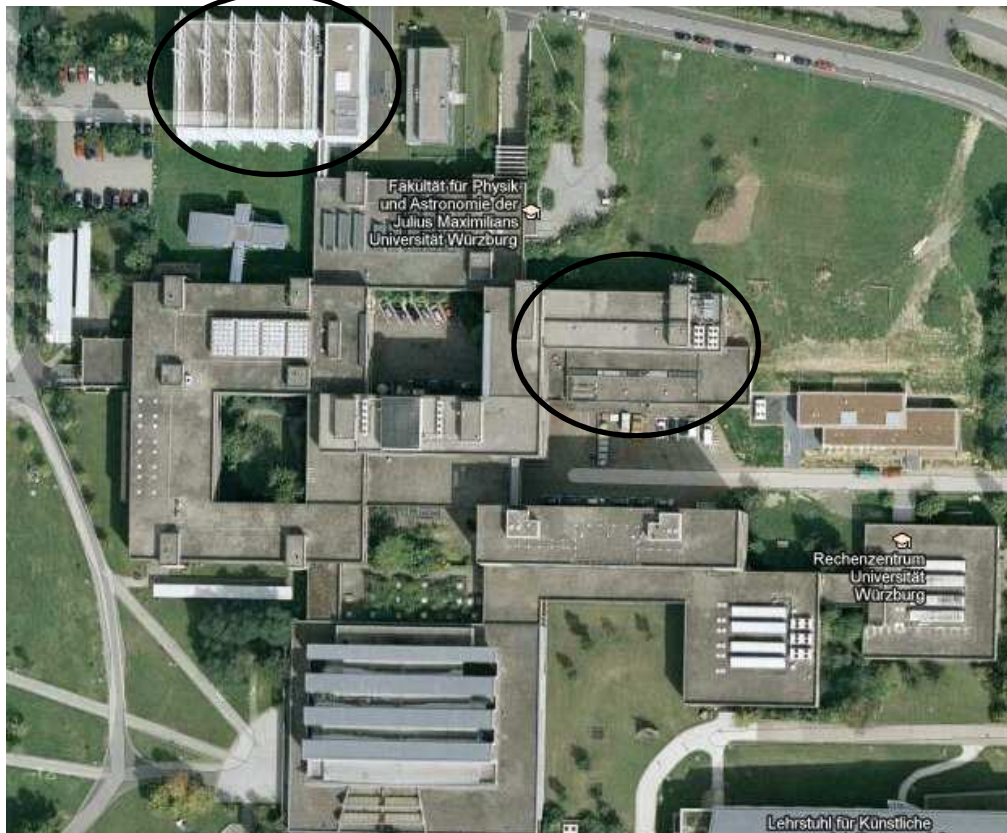
- Würzburg is located in the northern part of Bavaria.
- The region: **Franken**
- Population: 130,000
- Students: 24,300



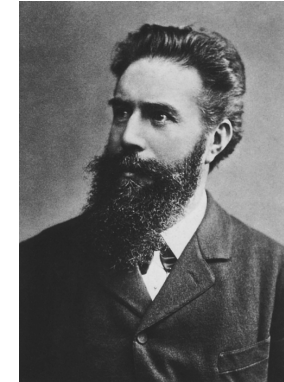
Würzburg is famous for:

- Wine (Wine producer of the year: Weingut Horst Sauer)
- And Dirk Nowitzki

Department of physics and astronomy



8 experimental physics chairs
5 theoretical physics chairs
+ several experimental and theoretical groups



Wilhelm C. Röntgen
Nobel Prize 1901
(X-rays)



Klaus von Klitzing
Nobel Prize 1985
(Quantum Hall effect)

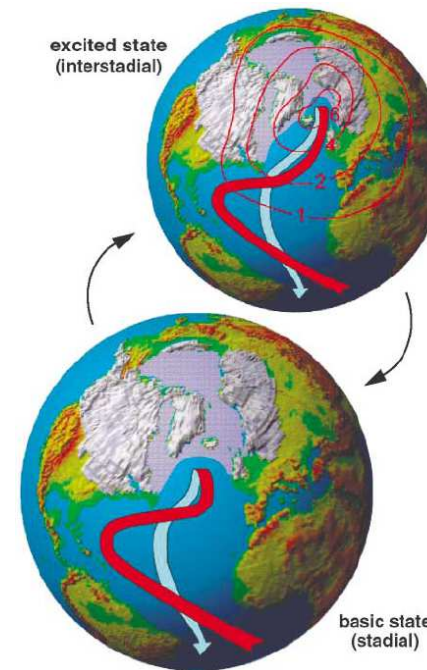
- Stochastic resonance: Weak signals can be enhanced by fluctuations (for a review Ref.[1])
- Ingredients:
 - Noise
 - Sub-threshold signal
 - Non-linear system, e.g. bistable systems
- SR as model was introduced to explain the periodic recurrences of ice ages: Benzi, Parisi, Sutera, Vulpiani [2]
- SR has been found in various systems, e.g. in crayfish mechanoreceptors [3]

Abrupt Glacial Climate Changes due to Stochastic Resonance

Andrey Ganopolski and Stefan Rahmstorf*

Potsdam Institute for Climate Impact Research, Box 601203, 14412 Potsdam, Germany

(Received 5 July 2001; published 4 January 2002)

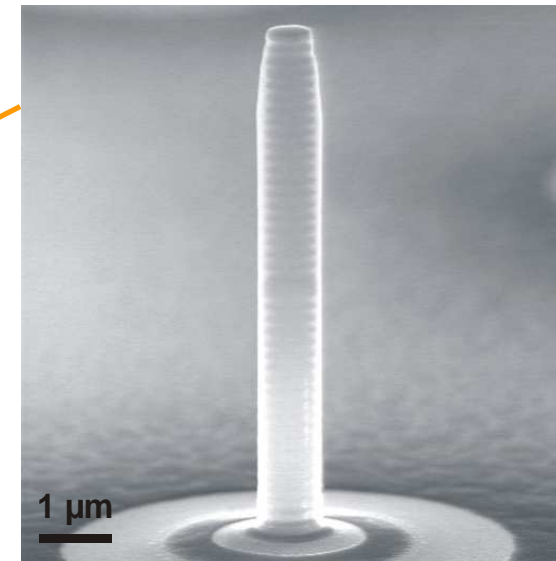
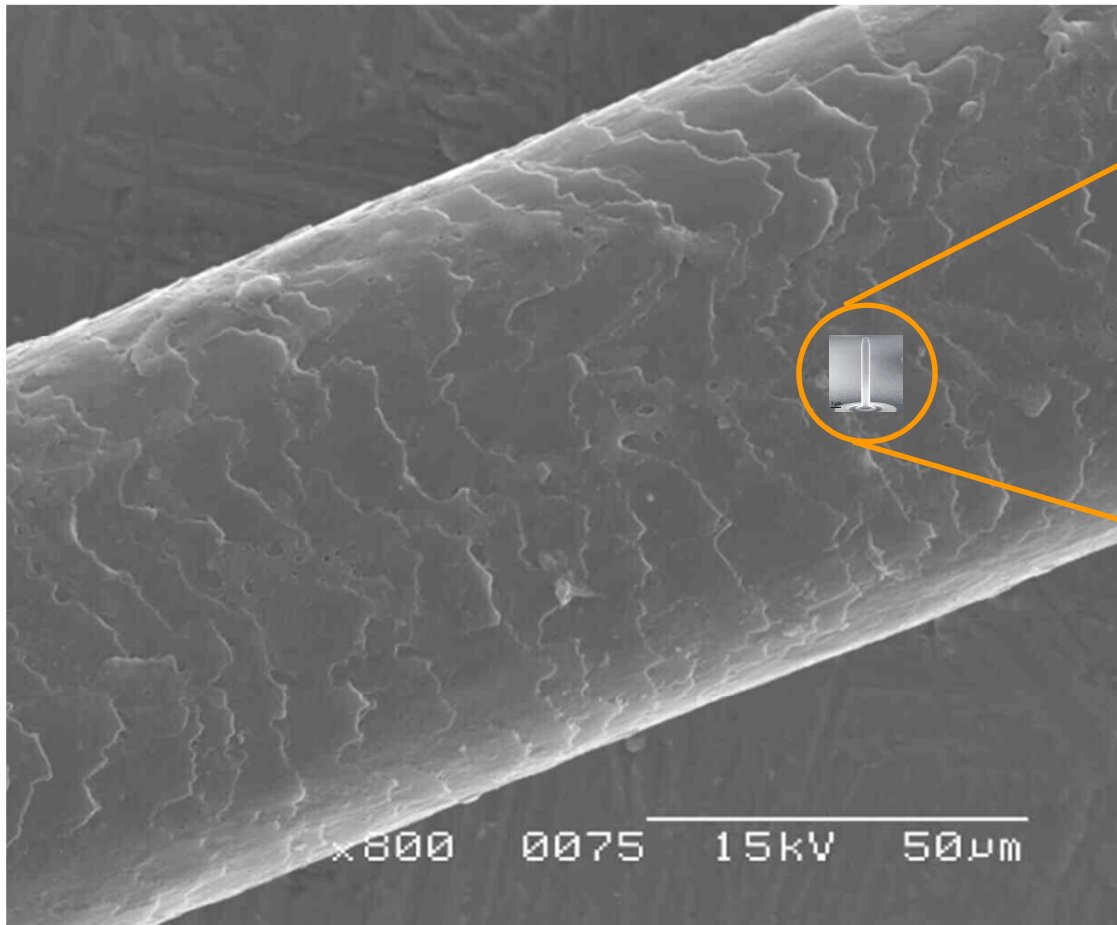


[1] L. Gammaitoni et al., "Stochastic resonance", *Reviews of Modern Physics*, Vol. 70, No. 1, January 1998

[2] Benzi, R., G. Parisi, A. Sutera, and A. Vulpiani, 1982, *Tellus* 34, 10.

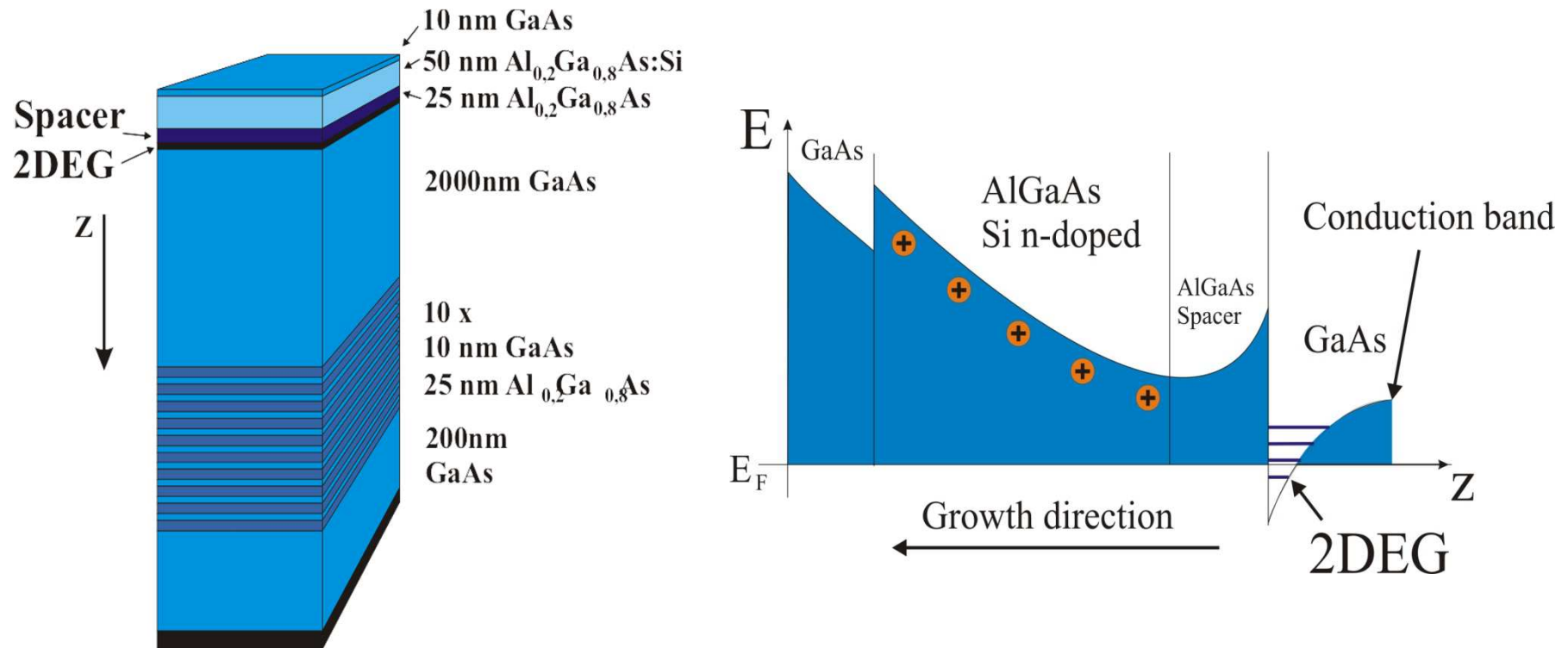
[3] Douglass, J. K., L. Wilkens, E. Pantazelou, and F. Moss, 1993, *Nature (London)* **365**, 337.

Electron microscope images of a human hair and a micro-pillar (fabricated @ our department)

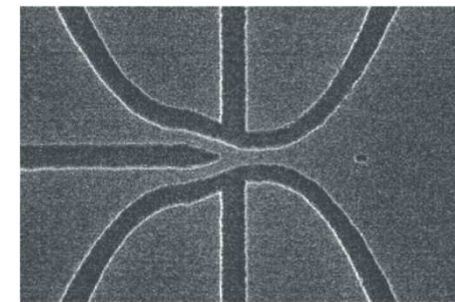
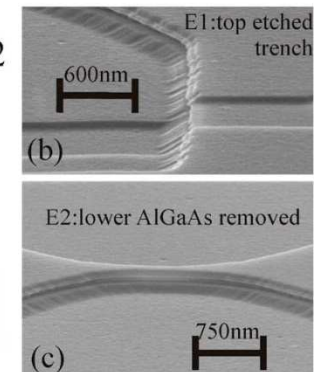
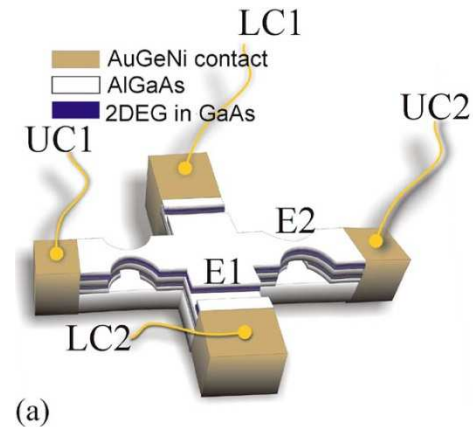
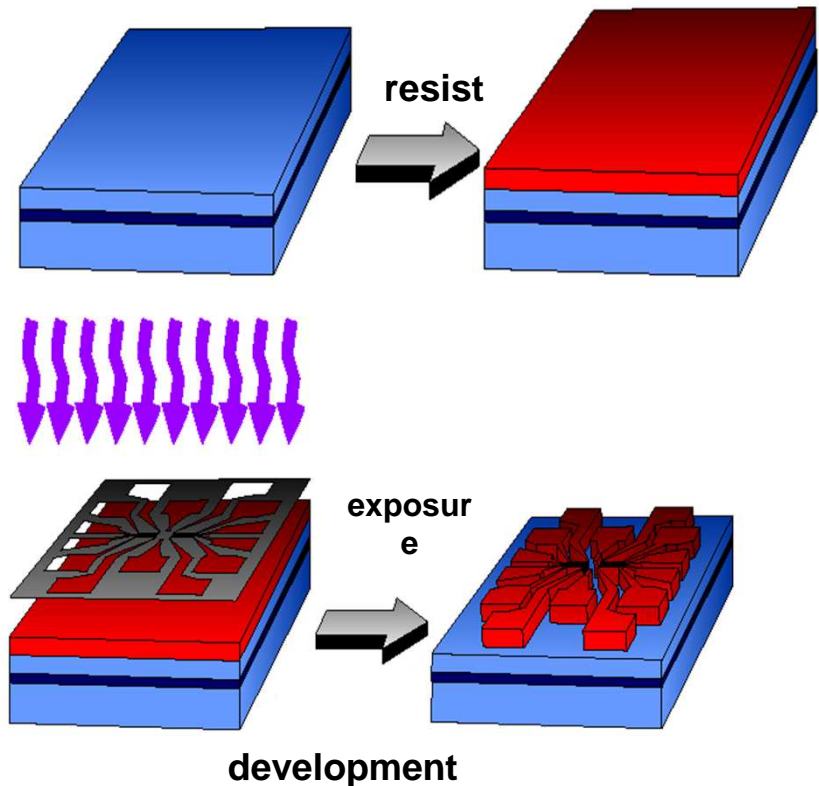


J.P. Reithmaier et al., „Strong coupling in a single quantum dot–semiconductor microcavity system“, *Nature* **432**, 197-200 (11 November 2004).

- **Growth, fabrication and transport properties of nanoelectronic devices**
 - Growth of 2DEGs and fabrication of electron waveguides.
 - Growth and fabrication of resonant tunneling diodes (RTD).
 - **Universal logic gate switching in resonant tunneling diodes (RTDs)**
 - Universal logic gate switching => NOR to NAND
 - Logic stochastic resonance (LSR)
 - **Stochastic resonance in nanoelectronic devices**
 - SR in electron waveguides
 - SR in RTDs for ac and periodic optical modulation
 - **Noise activated nonlinear dynamic sensors**
 - Magnetic field sensor based on (bistable) electron waveguides
 - **Energy harvesting: The quantum harvester class**
 - Transport as a consequence of state dependent diffusion.
 - Optimal energy to quanta conversion: A coupled QD system.
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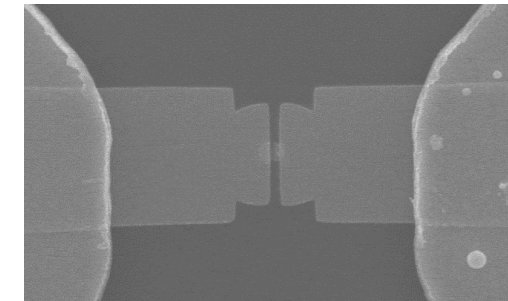
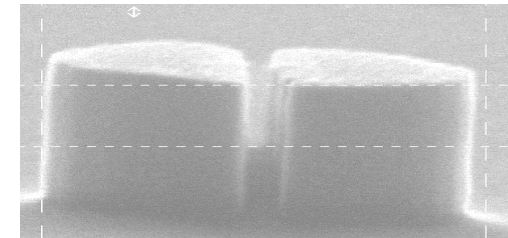
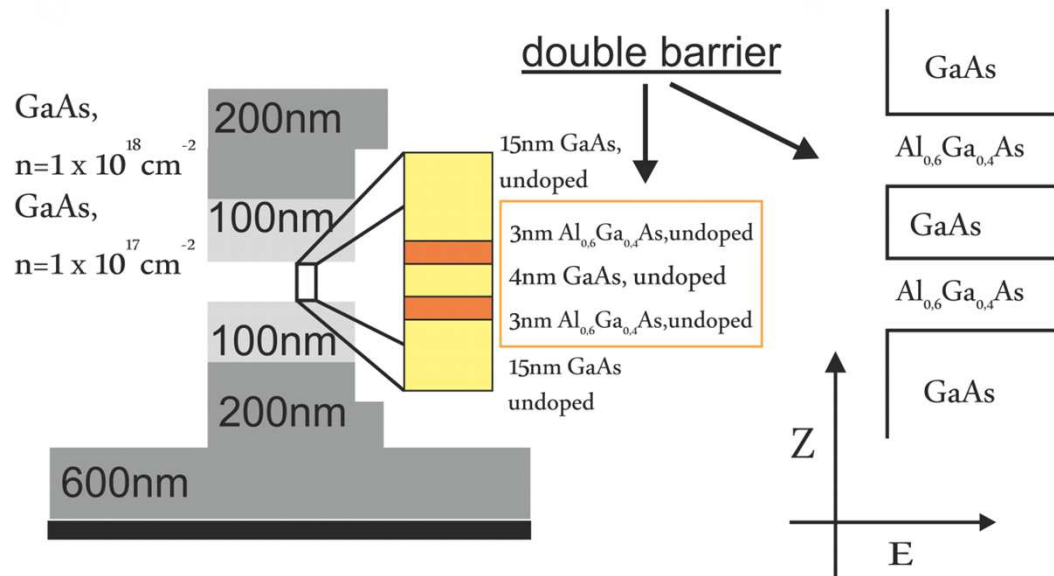
- Modulation doped GaAs/AlGaAs heterostructure (HEMT).
- Grown by molecular beam epitaxy.
- High mobility $\mu = 1.1 \cdot 10^6 \text{ cm}^2/\text{Vs}$ and charge density $n = 3.7 \cdot 10^{11} \text{ cm}^{-2}$



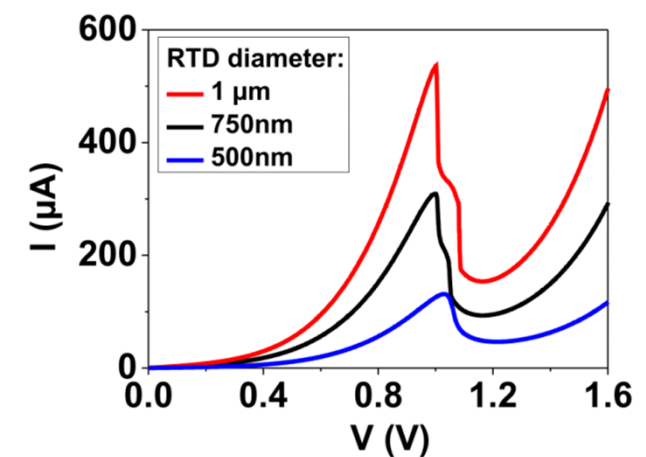
- Samples are grown by molecular beam epitaxy. →
- Electron beam or optical lithography. →
- Evaporating the etching mask (e.g. Cr) & Lift-off. →
- Remove the etching mask (HNO_3) →
- Resist (e.g. positive PMMA).
- Development of the resist.
- Wet or dry chemical etching (e.g. ECR-RIE)

DONE!! (plus contacts)

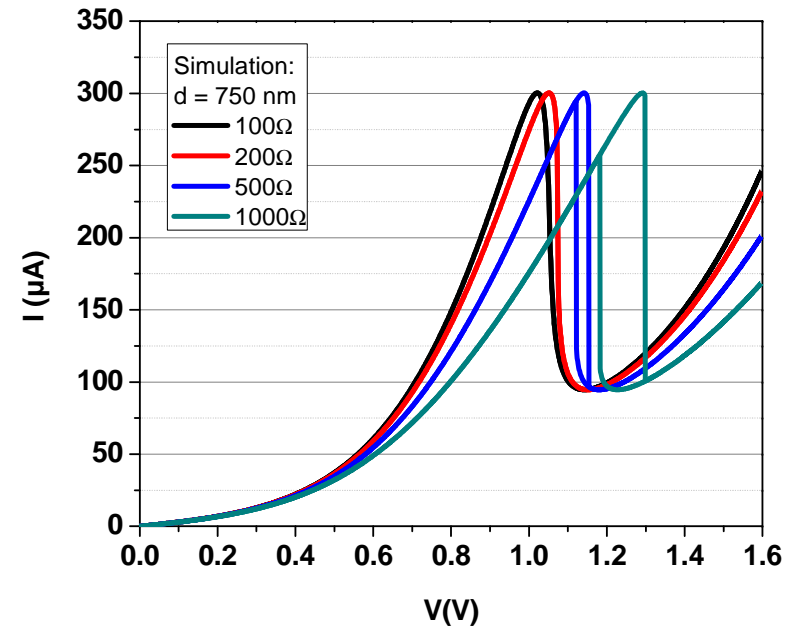
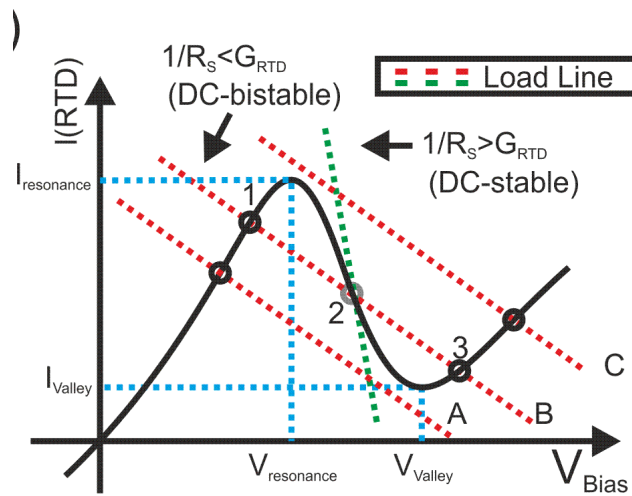
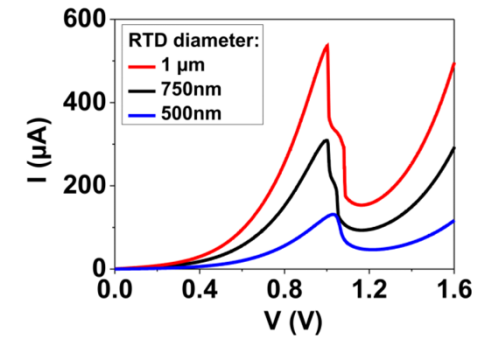
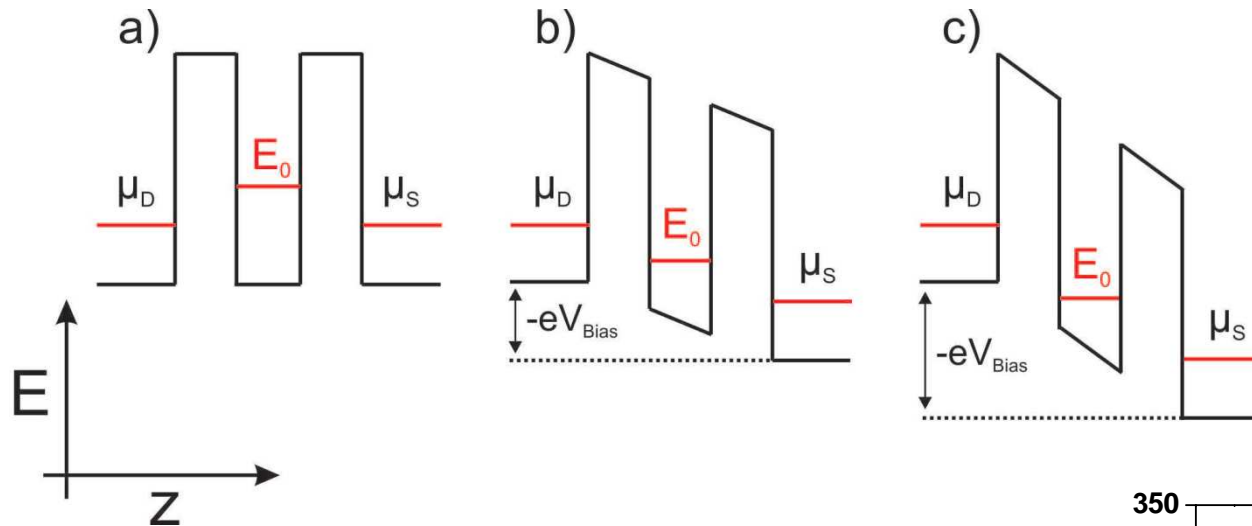
- RTDs based on the GaAs material system with AlGaAs/GaAs/AlGaAs double barriers.



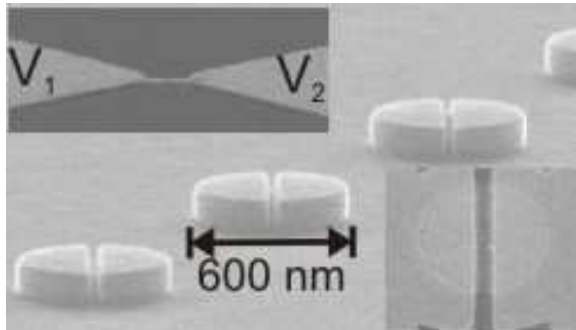
- Dry chemical etching is used to define RTD mesas from $12 \mu\text{m}$ down to 50 nm.
- BCB (polymer) for mesa isolation.
- Top Au/Ti/Ni contact.



Resonant tunneling diodes (RTDs): Tunable bistability via the load line effect

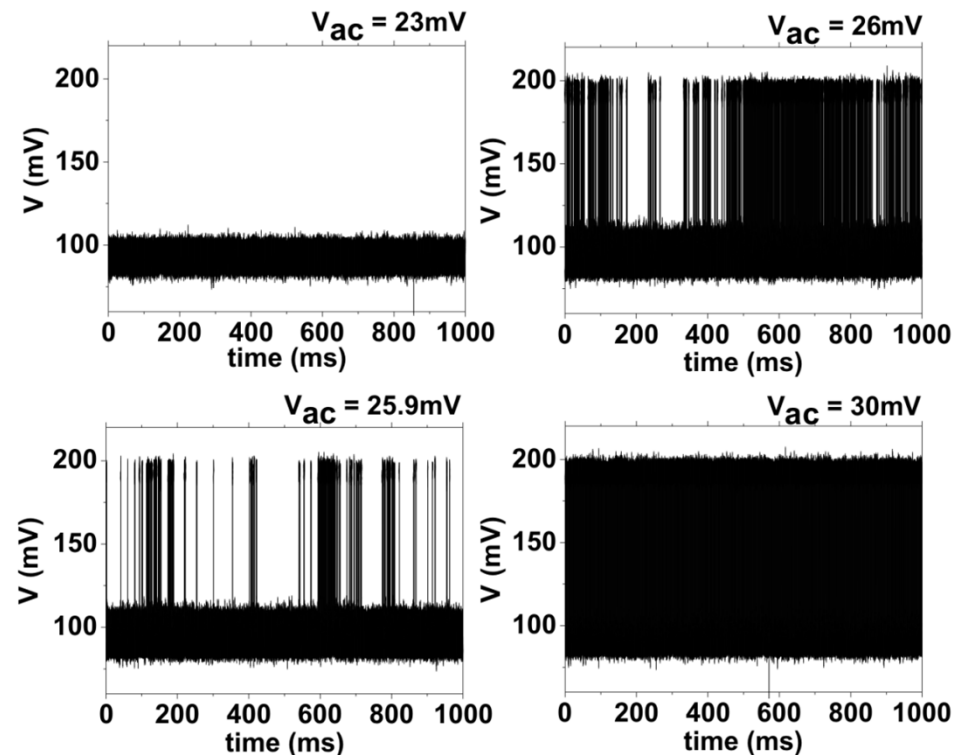


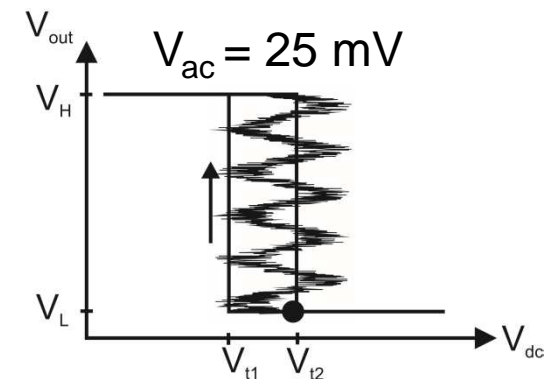
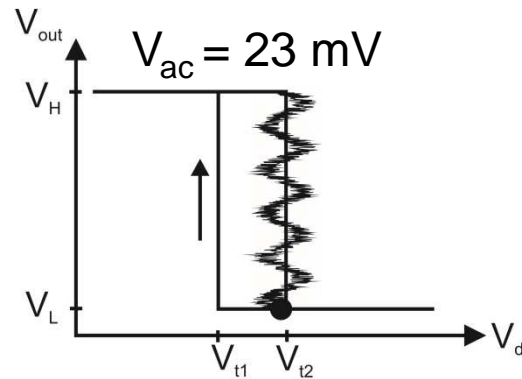
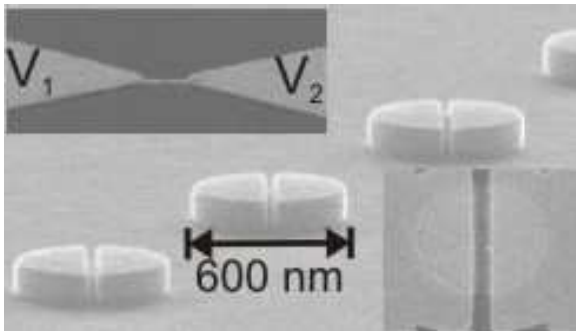
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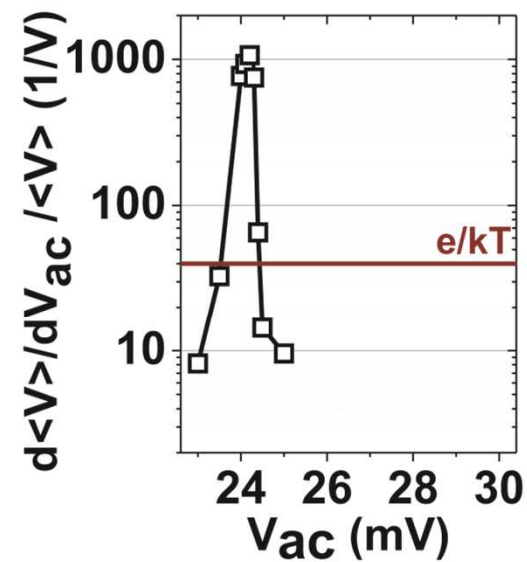
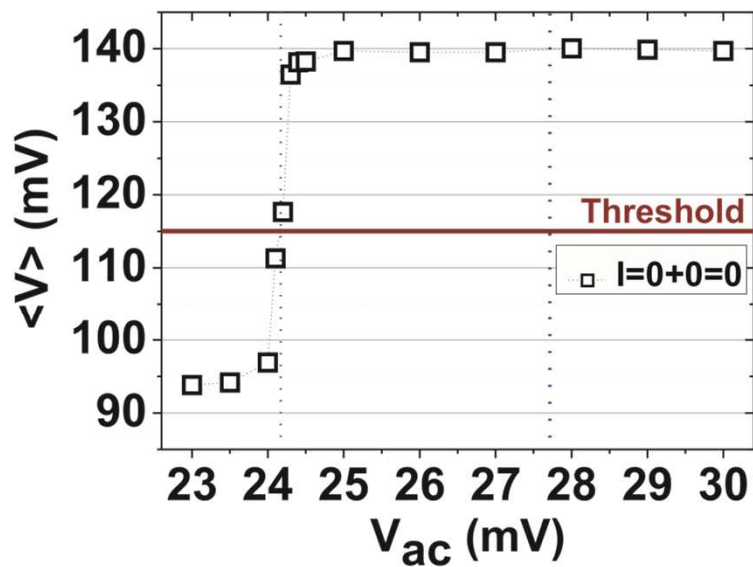
- Electron microscopy images of a trench etched RTD with diameter $d = 600$ nm
- Branches serve as logical inputs

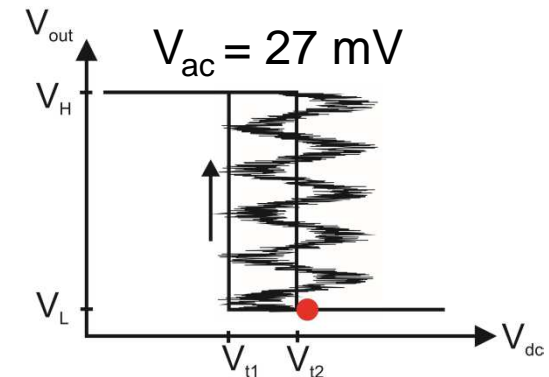
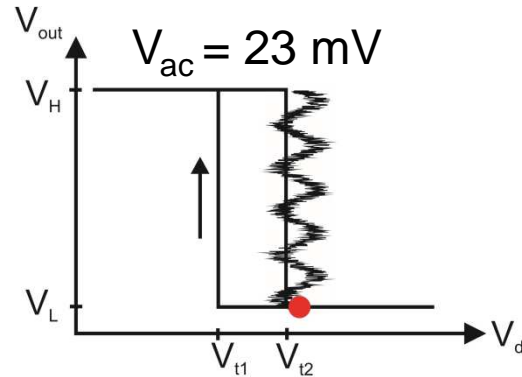
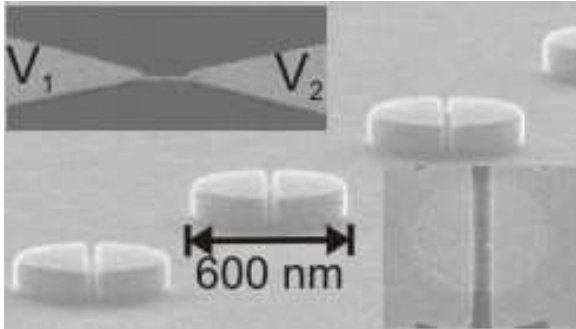
- Noise induced signal trains
- Mean value is efficiently controlled by input signals
- Can be integrated to arrays
- No classical kT limit of transconductance



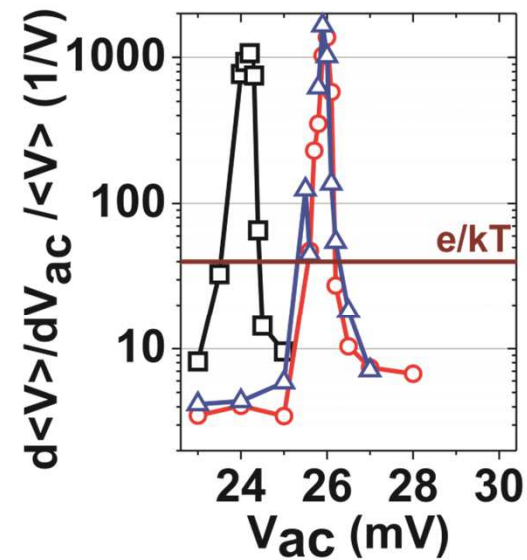
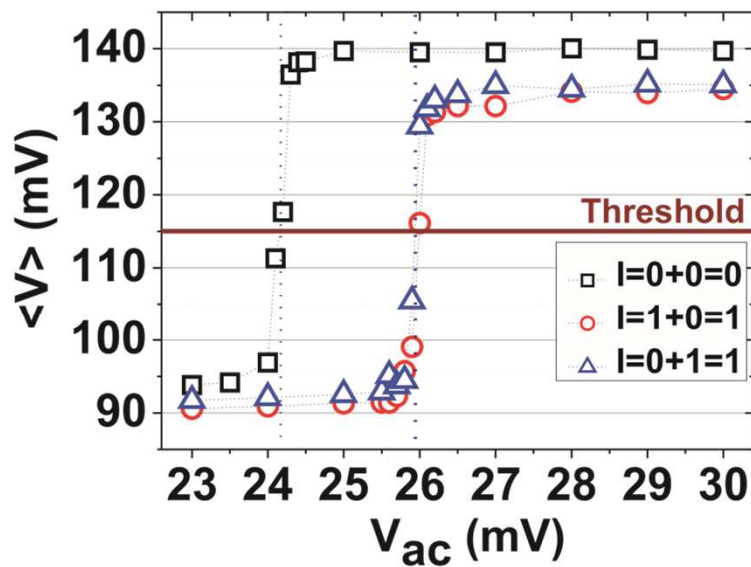


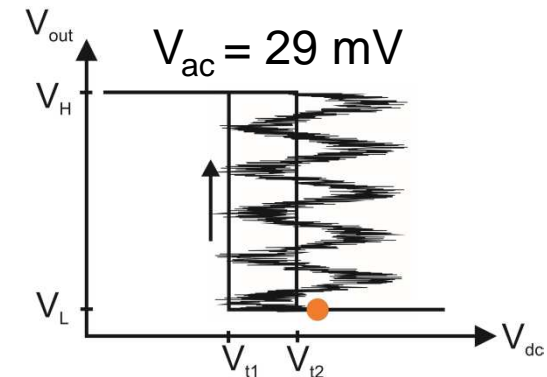
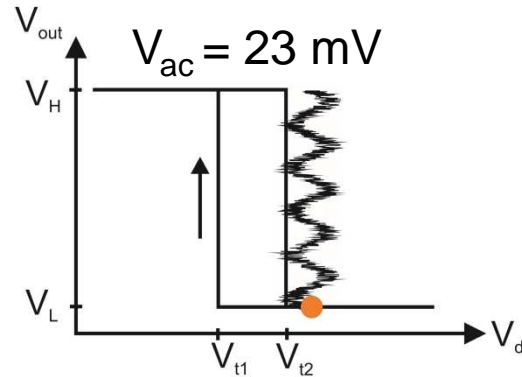
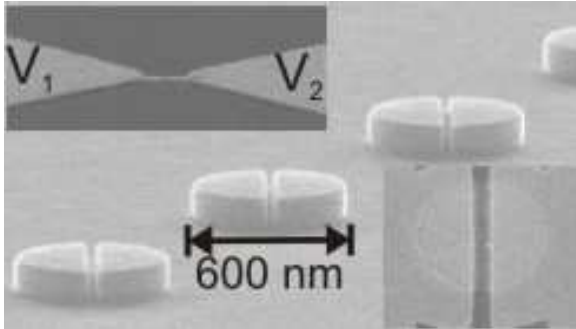
Switching voltages: $V_1 = V_2 = 0\text{mV}$



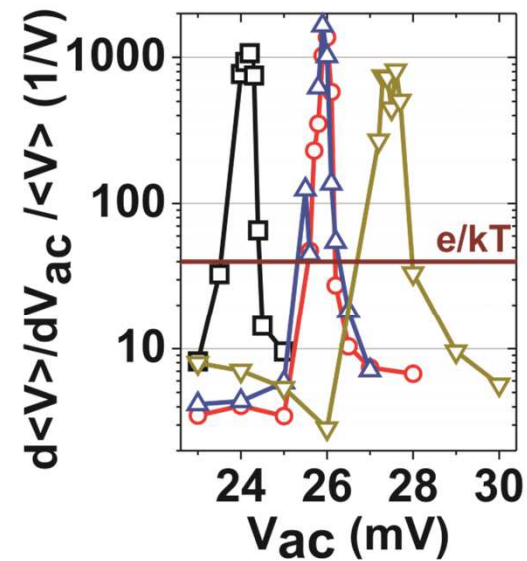
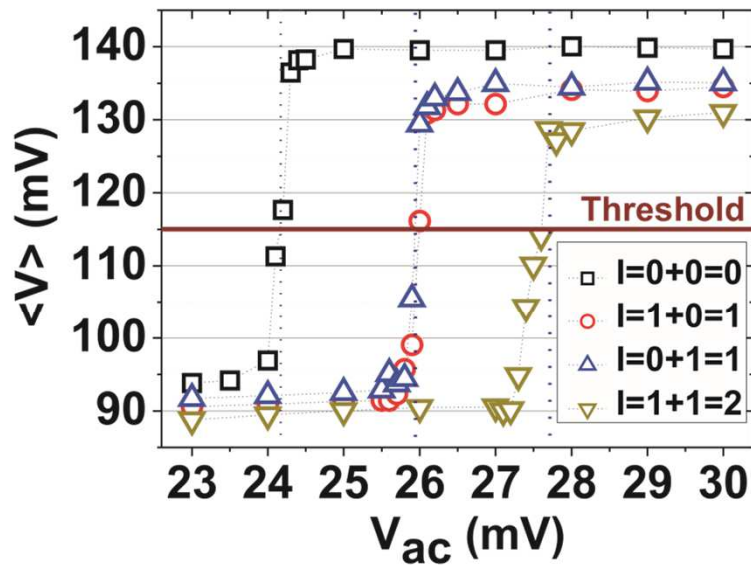


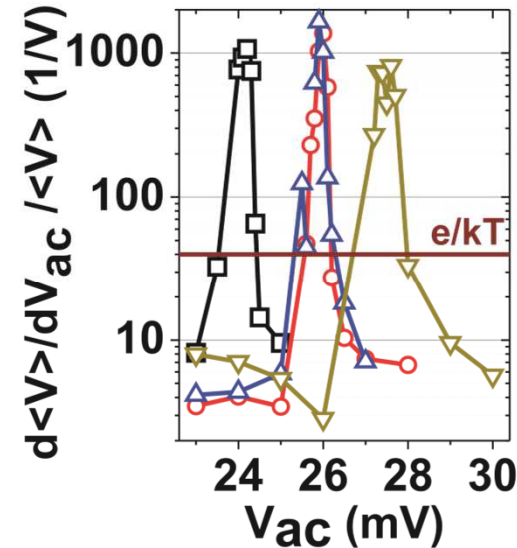
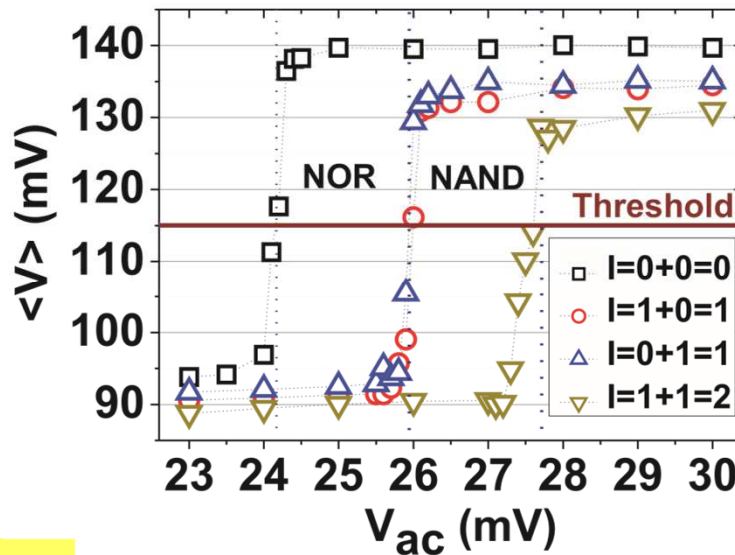
Switching voltages: $V_1 = 0, 2 \text{ mV}$ $V_2 = 2, 0 \text{ mV}$





Switching voltages: $V_1 = V_2 = 2 \text{ mV}$





NOR

| | | | |
|---|---|--|---|
| 0 | 0 | | 1 |
| 1 | 0 | | 0 |
| 0 | 1 | | 0 |
| 1 | 1 | | 0 |

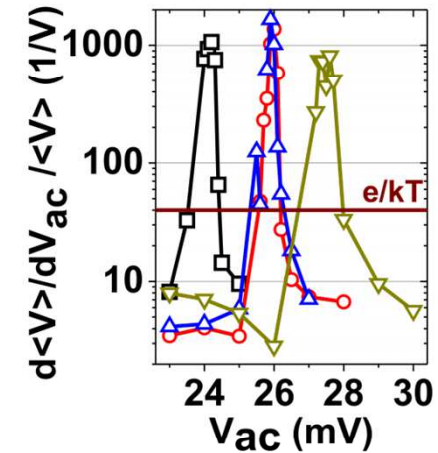
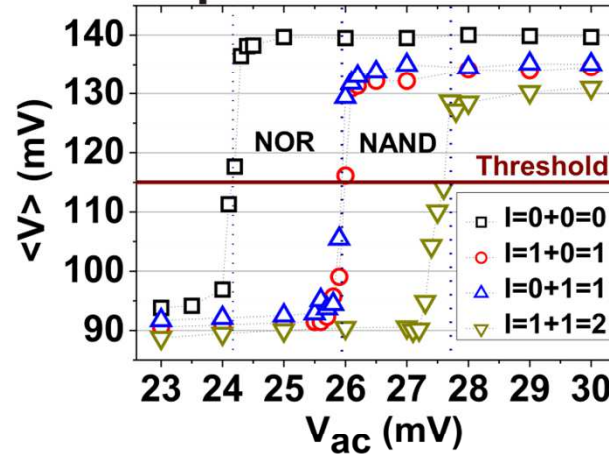


NAND

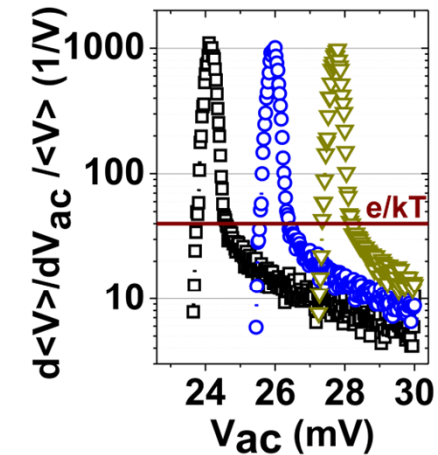
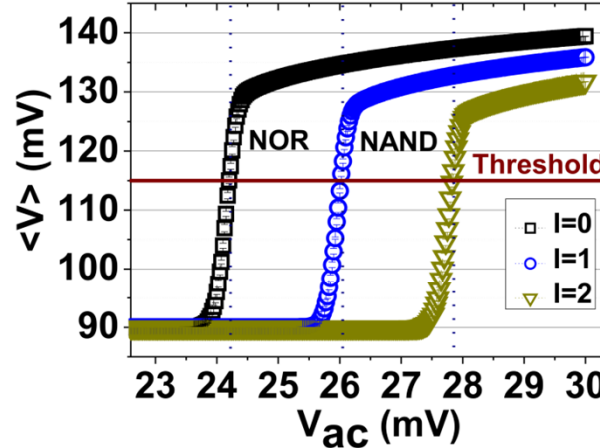
| | | | |
|---|---|--|---|
| 0 | 0 | | 1 |
| 1 | 0 | | 1 |
| 0 | 1 | | 1 |
| 1 | 1 | | 0 |

- Switch from NOR to NAND for $\Delta V_{ac} < 1$ mV with a logic input voltage 2 mV.

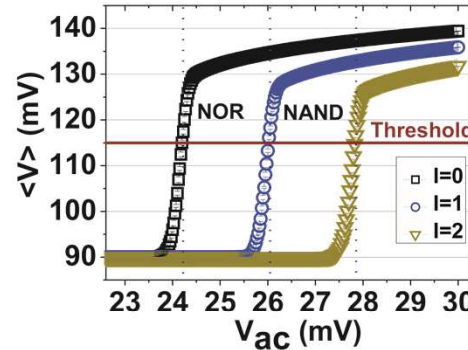
experiment



simulation

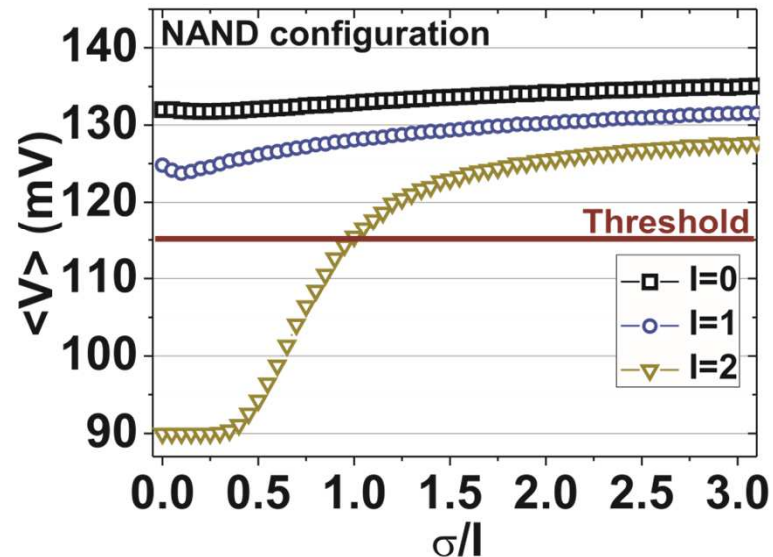
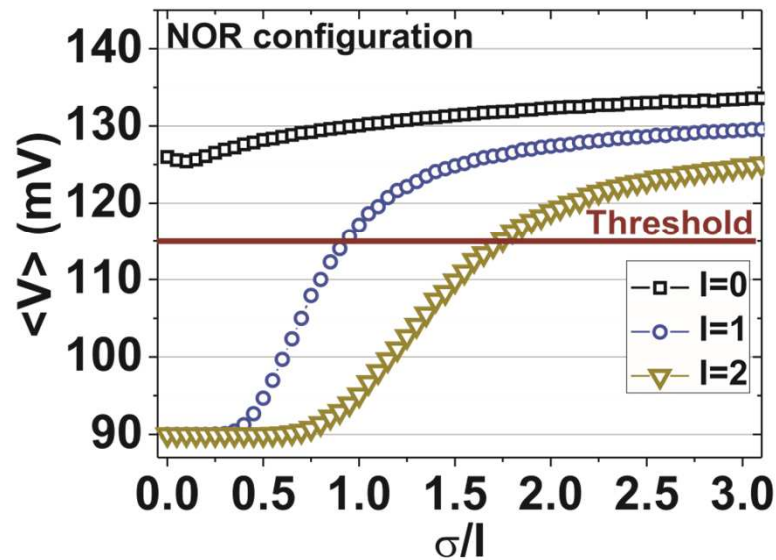


- ☐ Schmitt-Trigger simulation.
- ☐ All Parameters from the experiment.
- ☐ Excellent agreement !!

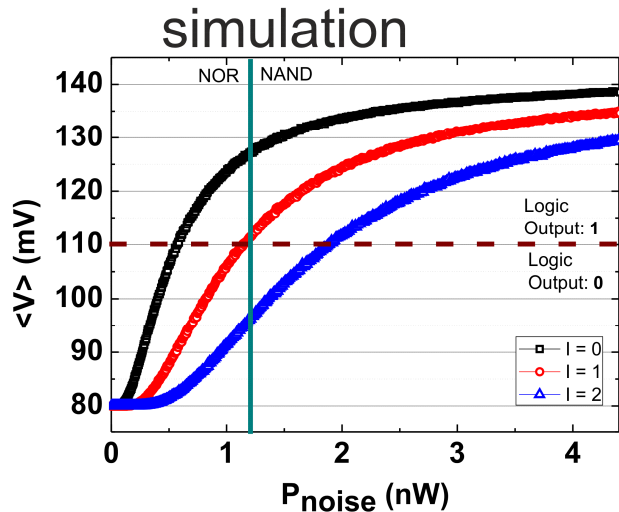
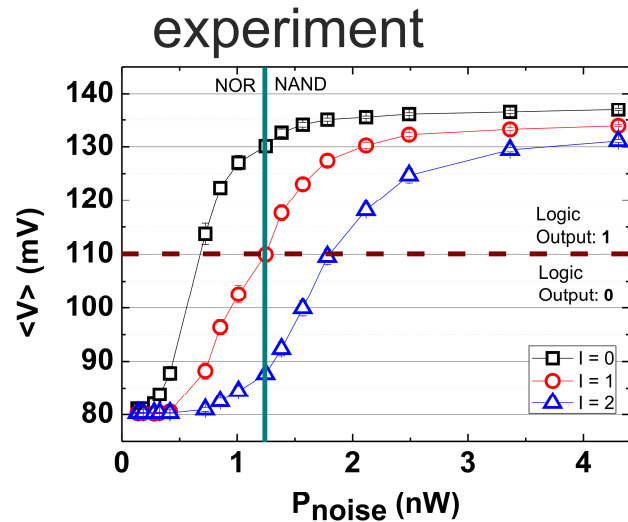


$V_{ac} = 24.6 \text{ mV}$

$V_{ac} = 26.6 \text{ mV}$



□ Robust response to the noise floor up to **100%** of logic input

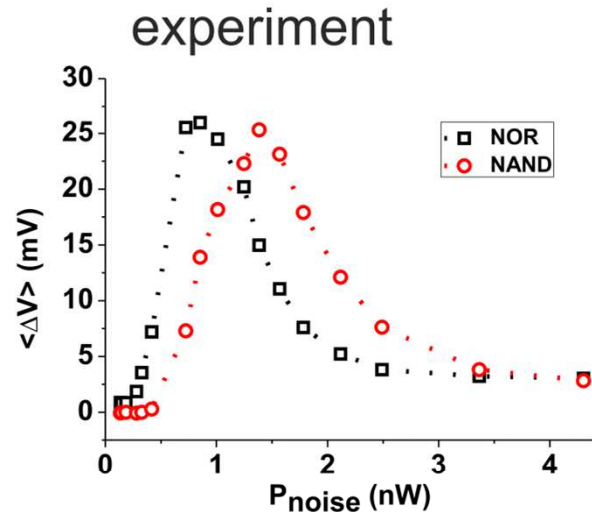


Previous:

- Universal logic gate switching controlled by the amplitude of the periodic forcing V_{ac} .

Now:

- Universal logic gate switching solely controlled by the noise floor.
- Two universal logic gates: NOR/NAND.

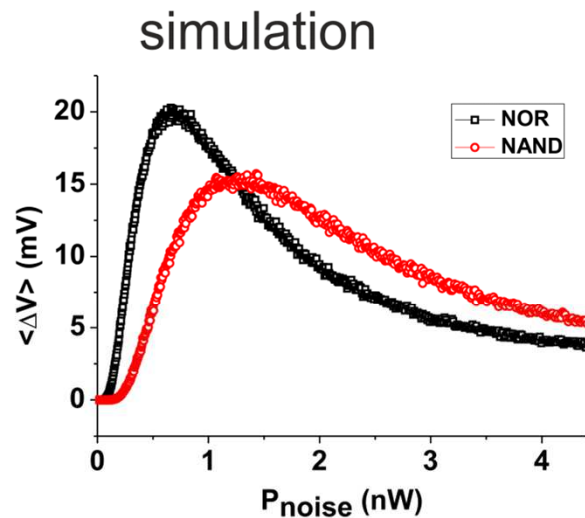


For the logic NOR gate:

- The mean value difference is defined as

$$\langle V \rangle = V(I=0) - V(I=1)$$

$P_{\text{noise}} = 0.9$ nW the maximum corresponds to the logic NOR



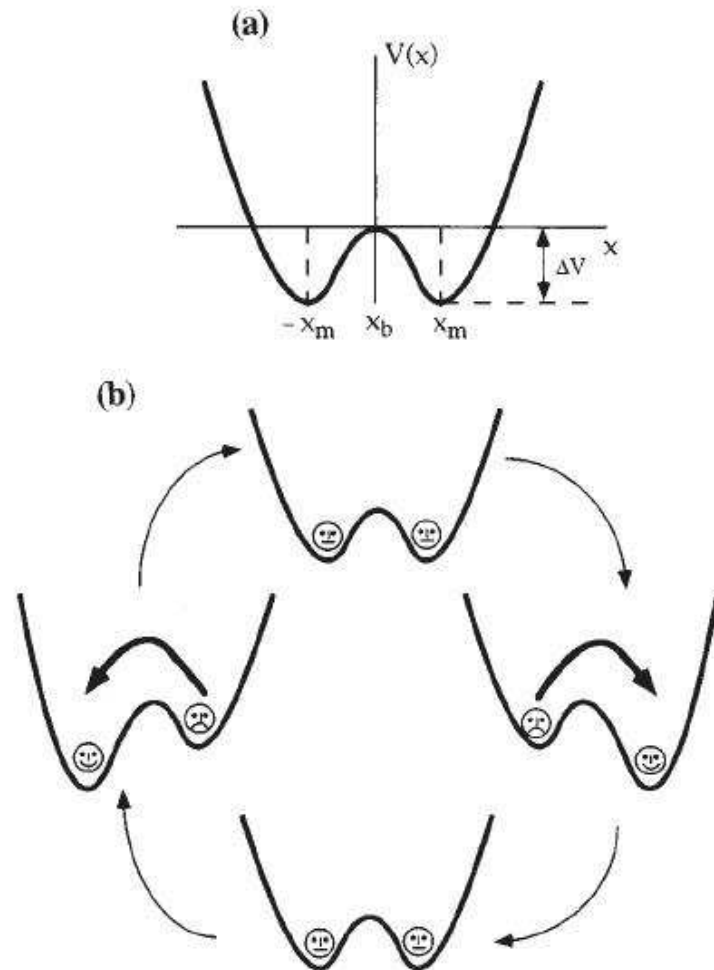
For the logic NAND gate:

- The mean value difference is defined as

$$\langle V \rangle = V(I=1) - V(I=2)$$

$P_{\text{noise}} = 1.4$ nW the maximum corresponds to the logic NAND

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Overdamped motion of a Brownian particle in a bistable potential in the presence of noise and periodic forcing

$$\dot{x} = -V'(x) + A_0 \cos(\omega t + \varphi) + \xi(t)$$

with

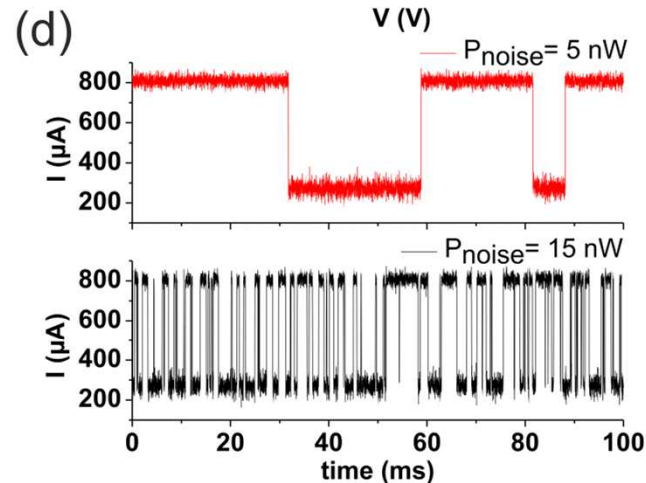
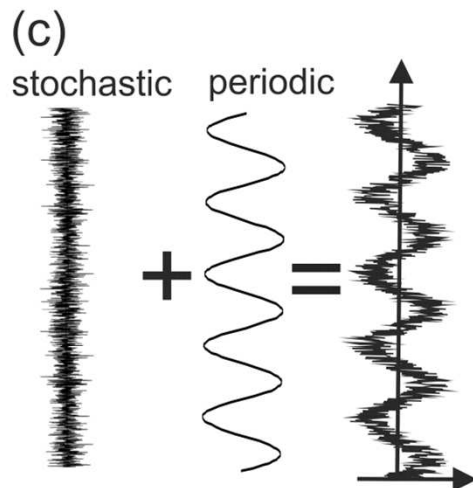
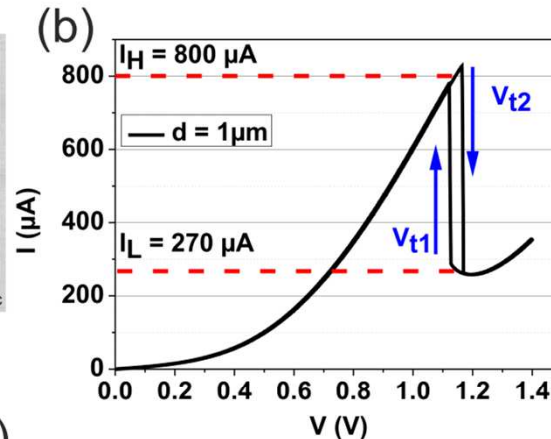
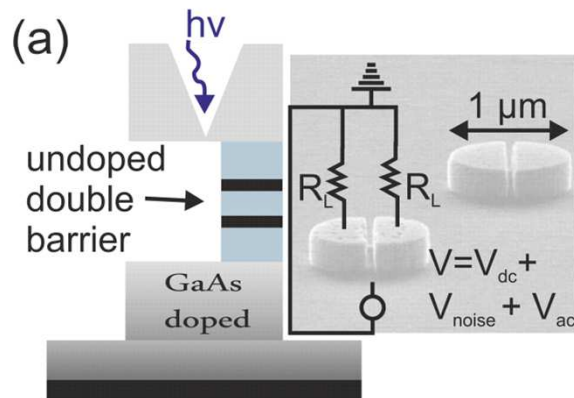
$$V(x) = -\frac{1}{2}x^2 + \frac{1}{4}x^4$$

Noise-induced hopping between the local equilibrium states with the Kramers rate

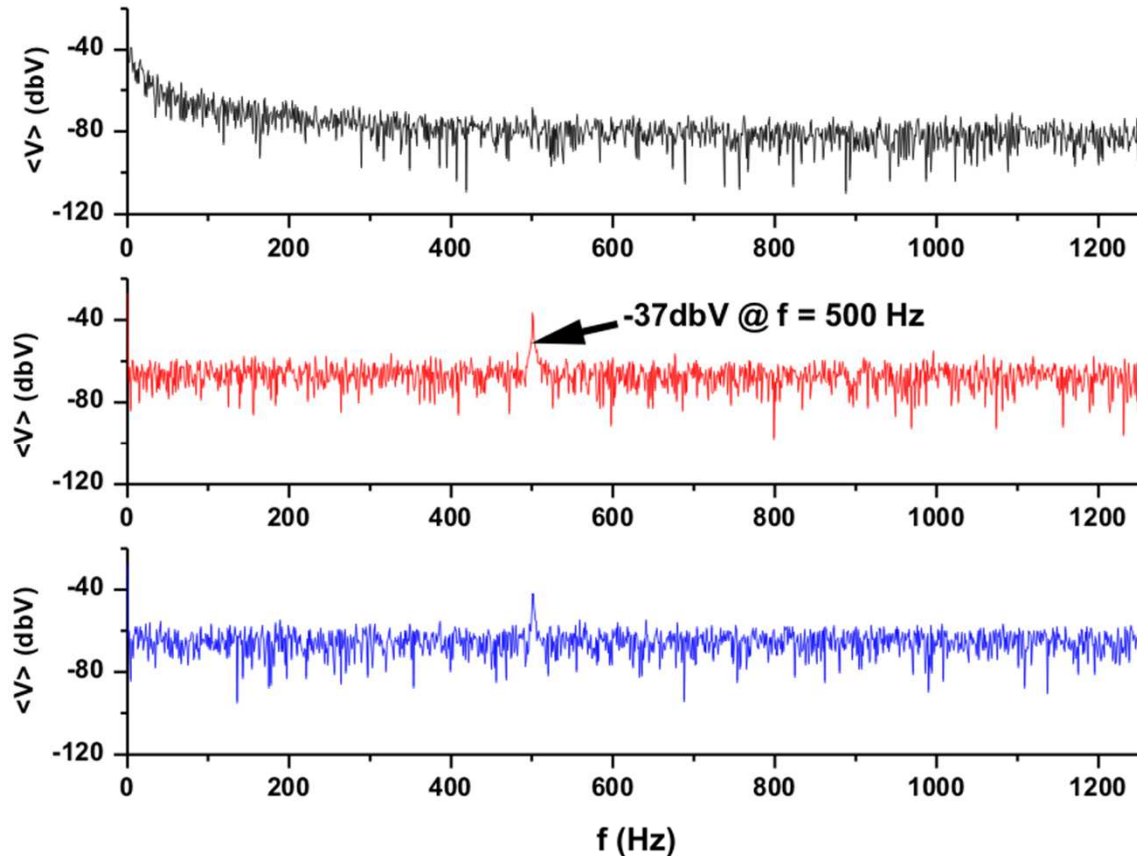
$$r_K = \frac{1}{\pi\sqrt{2}} \exp\left(-\frac{\Delta V}{D}\right)$$

The *time-scale matching condition* for stochastic resonance:

$$T_\omega = 2T_K$$



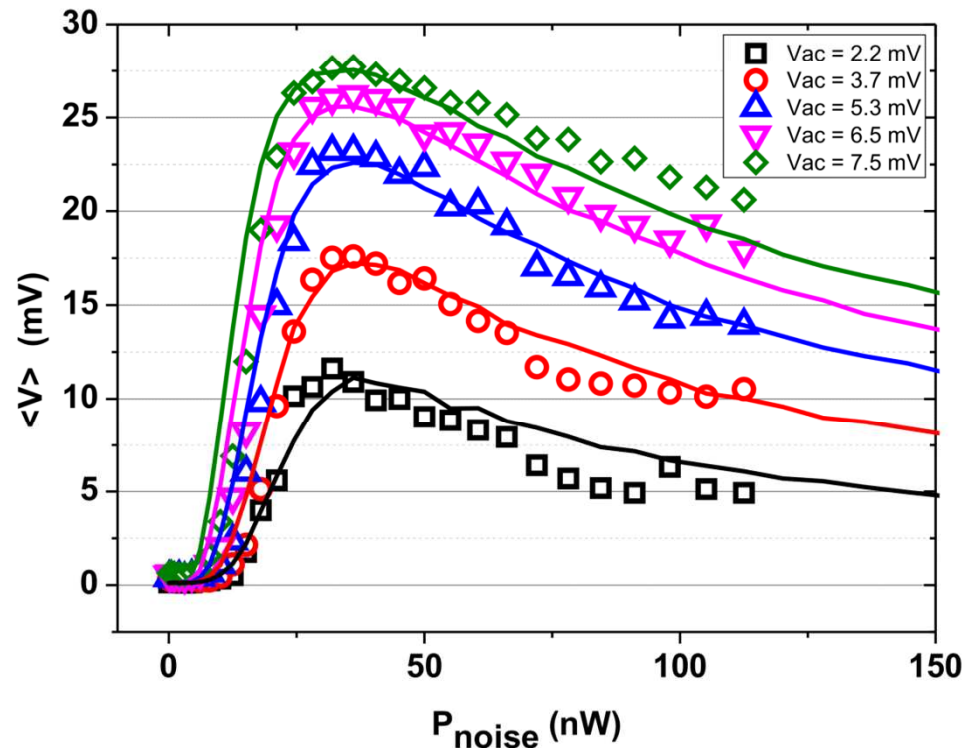
- RTD is bistable with stable outputs $I_H = 800 \mu\text{A}$ and $I_L = 270 \mu\text{A}$.
- Works @ RT
- PVR ~ 3
- Noise induced switching between the two stable states appear.
- Time scale T_K is given by the inverse of the Kramer's rate.



- For $P_{\text{noise}} < P_{\text{SR}}$ no spectral component at $f = 500$ Hz is found.

- For $P_{\text{noise}} > P_{\text{SR}}$ the spectral component at $f = 500$ Hz is still apparent.

- At the optimum noise level P_{SR} , the spectral amplitude reaches a maximum value and is decreasing apart from P_{SR} .



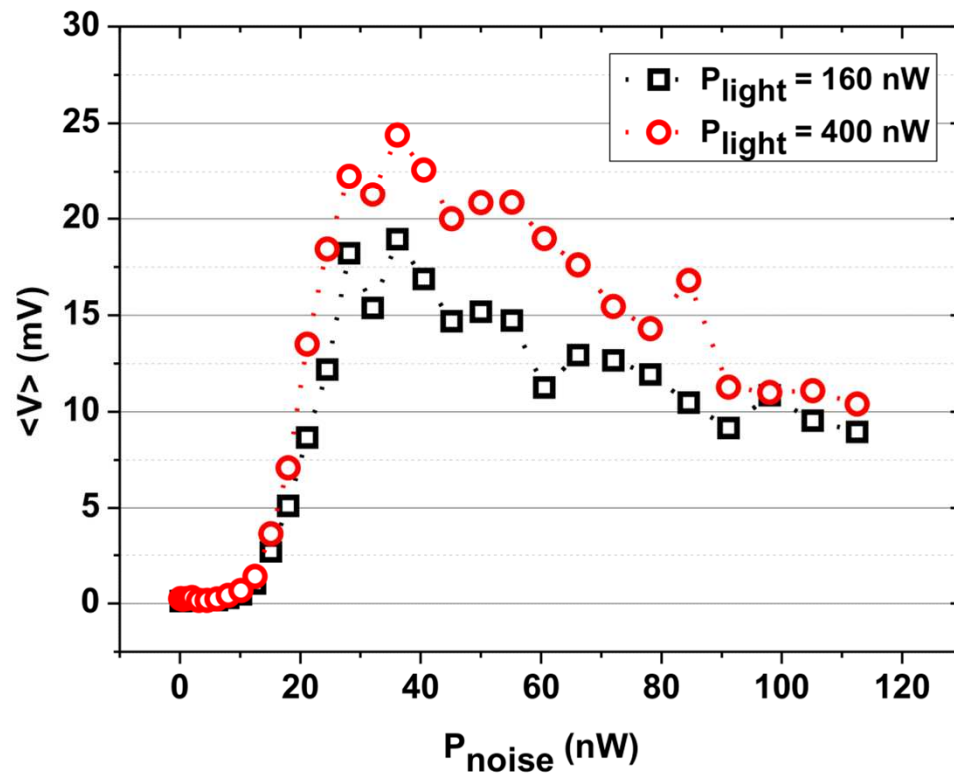
- For $P_{\text{noise}} < P_{\text{SR}}$ the spectral component at $f = 500 \text{ Hz}$ is increasing.
- Maximum synchronization @ $P_{\text{SR}} \Rightarrow \text{SR}$.
- For $P_{\text{noise}} > P_{\text{SR}}$ the spectral component is decreasing again.

Simulations (solid):

- Ideal two state model (Schmitt Trigger) with parameters from the experiment.
- e.g. the barrier height was set to 16 mV as the hysteresis width of the device was 32 mV.

Now:

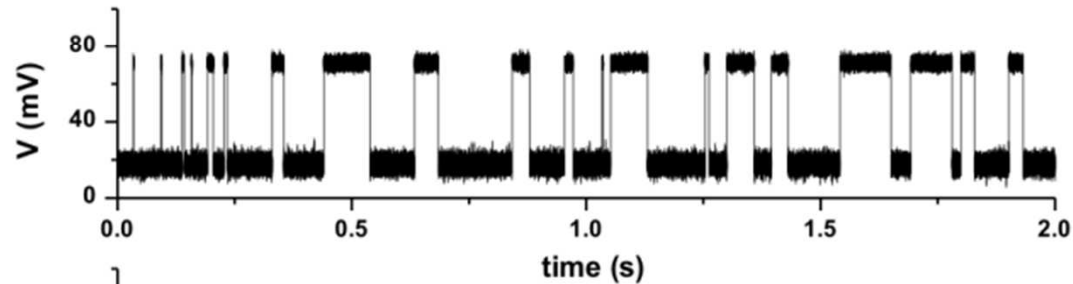
- Change from ac modulation to a periodic light modulation.
- Energy of the light $E = 2.73\text{ eV}$ (448nm) above the GaAs bandgap.
- Mechanically chopped light signal at $f = 500\text{ Hz}$.



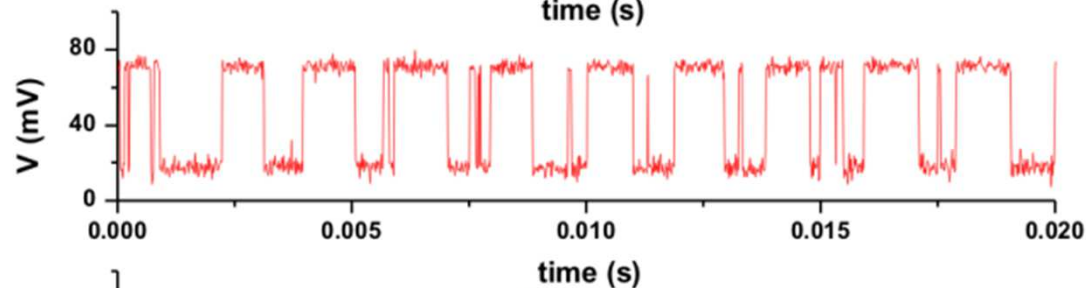
- For $P_{\text{noise}} < P_{\text{SR}}$ the spectral component at $f = 500\text{ Hz}$ is increasing.

- Maximum synchronization @ P_{SR}
=> **SR**.

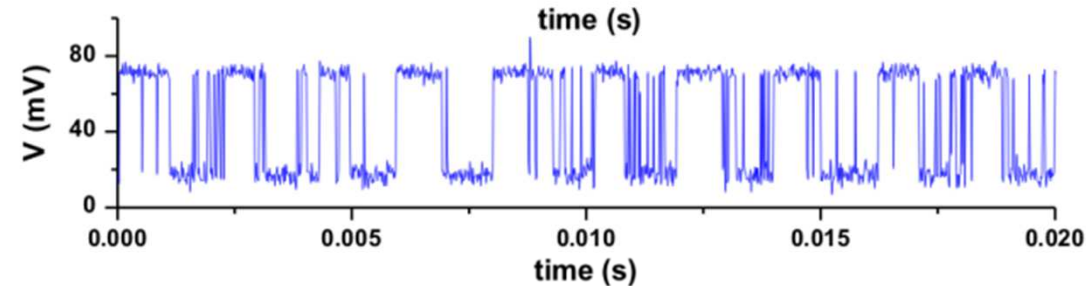
- For $P_{\text{noise}} > P_{\text{SR}}$ the spectral component is decreasing again.



$$P_{\text{noise}} = 2 \text{ nW}$$



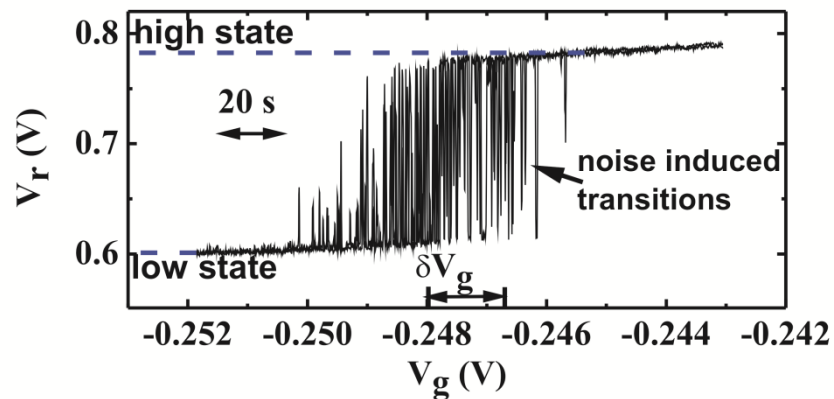
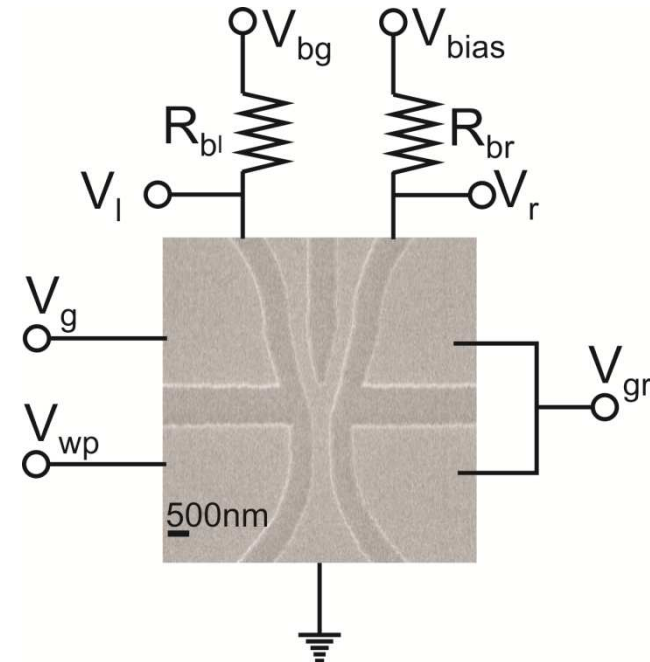
$$P_{\text{noise}} = 32 \text{ nW}$$



$$P_{\text{noise}} = 112 \text{ nW}$$

At $P_{\text{noise}} = 32 \text{ nW}$ the output follows almost perfectly the input signal !!

- The input and the working point voltages set the condition of the Y-branch switch.
- Self-gating leads to a bistable transfer characteristic.
- Noise induced oscillations occur
- All measurements @ 20K.



Input signal:

$$V_g(t) = V_{g,0} + \delta V_g \cdot \sin(\omega t)$$

Weak periodic signal:

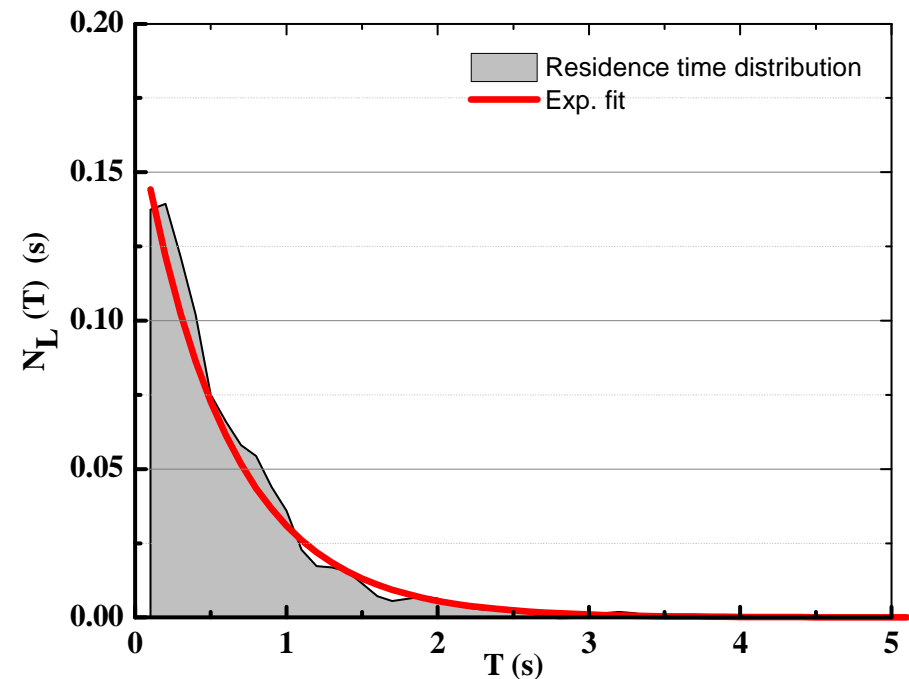
$$\delta V_g = 1.3 \text{ mV}$$

- For the unmodulated system, e.g. $f = 0$ Hz, the residence time distribution decays exponentially.
- The exponential decay is the inverse of the Kramer's rate and given by T_k :

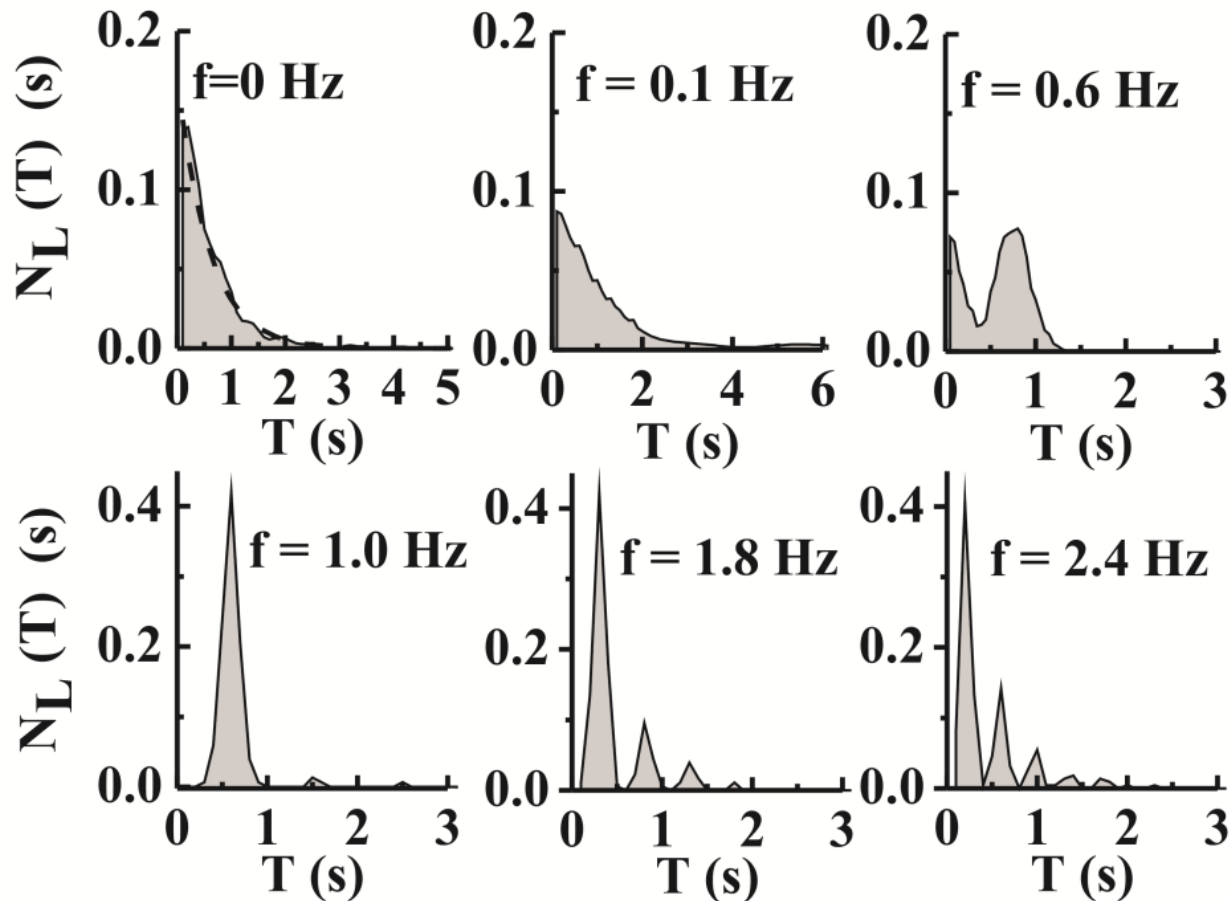
$$N_{L,H}(T) \propto \exp\left(-\frac{T}{T_K}\right)$$

From fitting:

$$T_K = (0.502 \pm 0.044)s$$



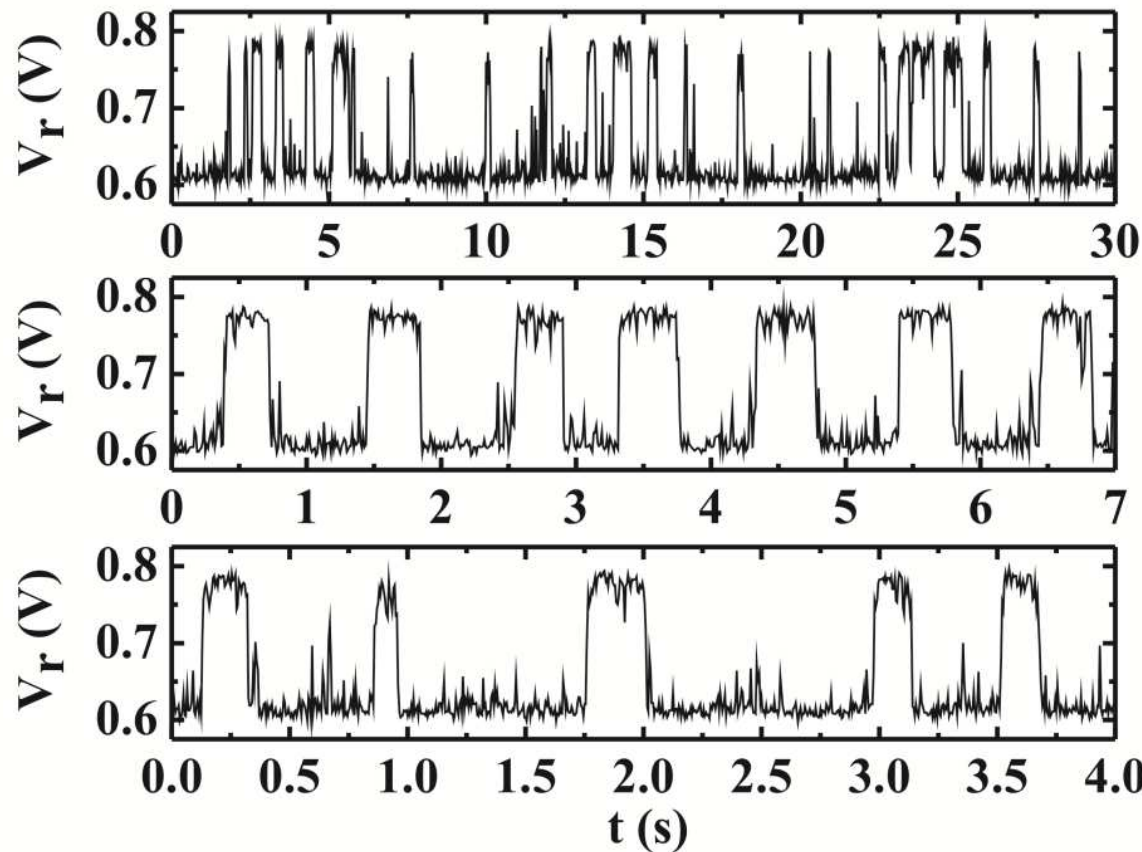
Time matching condition of SR: $T_\omega = 2T_K$



- For $f < f_{SR}$ the residence time distribution is strongly controlled by the noise
- For $f > f_{SR}$ odd multiples of the periodic forcing T_ω occur:

$$T_n = (2n - 1)T_\omega / 2$$

At the optimum frequency $f = 1$ Hz the residence time distribution is almost perfectly restricted to the first peak.



$f = 0.1$ Hz

$f = 1$ Hz

$f = 1.8$ Hz

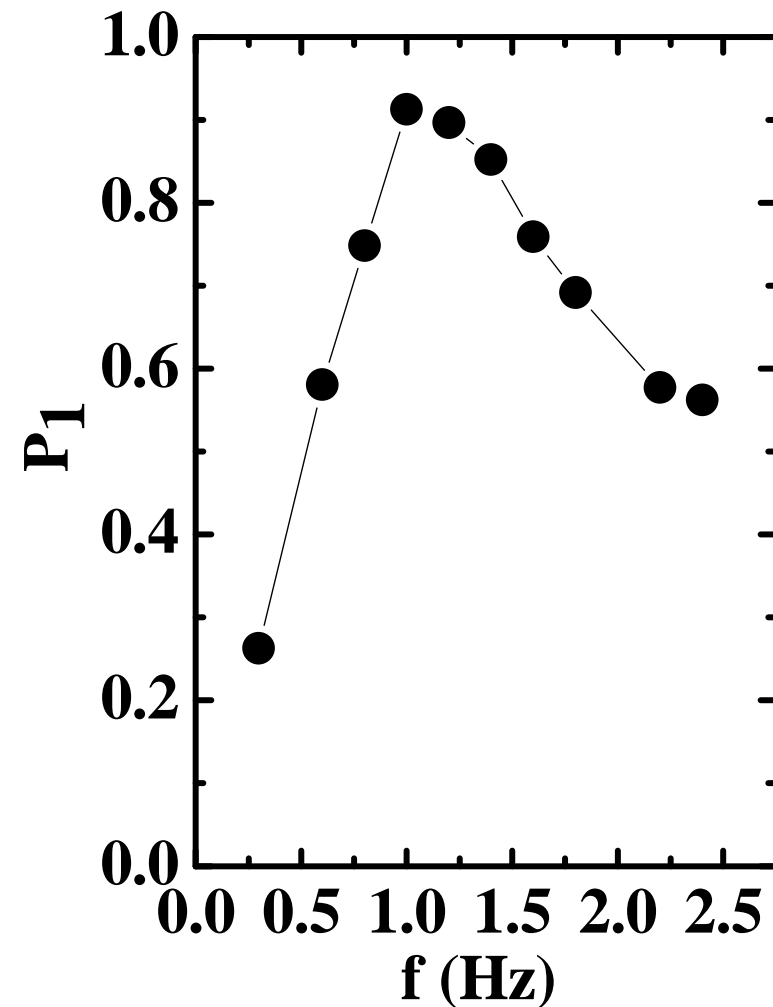
At $f = 1$ Hz the noise dynamics and the external (weak) periodic forcing are synchronized => Stochastic resonance.

- The strength P_1 of the first peak at $T_\omega/2$ (the area under the peak) is a measure of the synchronization between the periodic forcing and the switching between the wells.
- P_1 is defined as

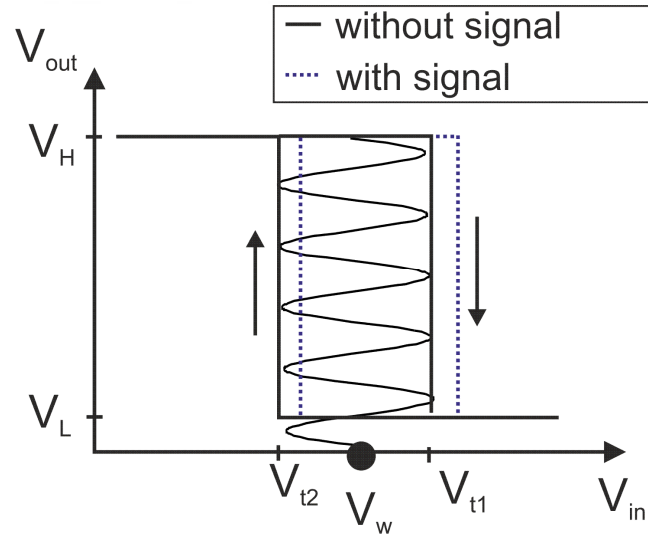
$$P_1 = \int_{T_1 - \alpha T_\omega}^{T_1 + \alpha T_\omega} N_L(T) dT$$

With $n=1,2,\dots$

And $0 < \alpha = 0.2 < 0.25$



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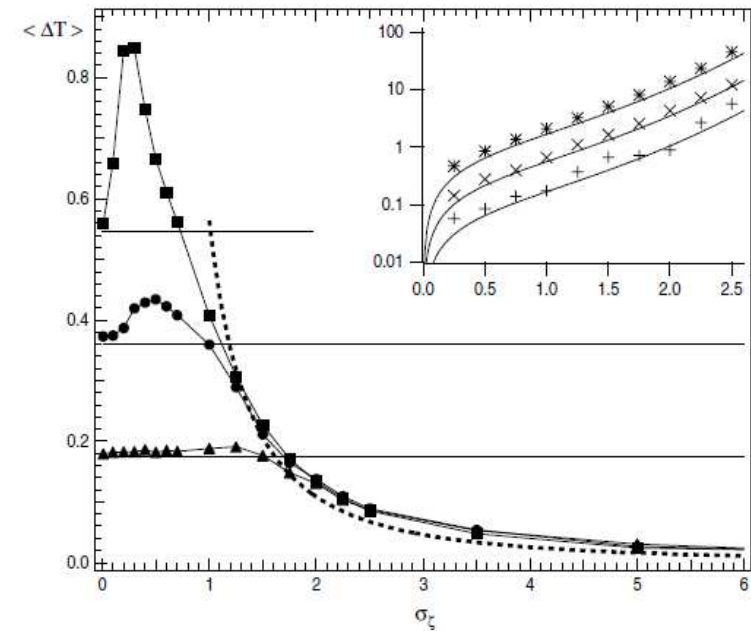


The variable of interest is the residence time difference ΔT between the time spend in the two stable states $T_{H,L}$ with

$$\Delta T = T_H - T_L$$

The response of $\langle \Delta T \rangle$ for large noise intensity σ_ξ is (expanded to first order in ϵ (target signal)):

$$\langle \Delta T \rangle = 4\epsilon \sqrt{\pi \tau / \sigma^2} \exp[b^2 / 2\sigma_\xi^2] \operatorname{erf}\left(\frac{b}{\sqrt{2\sigma_\xi^2}}\right) + O(\epsilon^2)$$

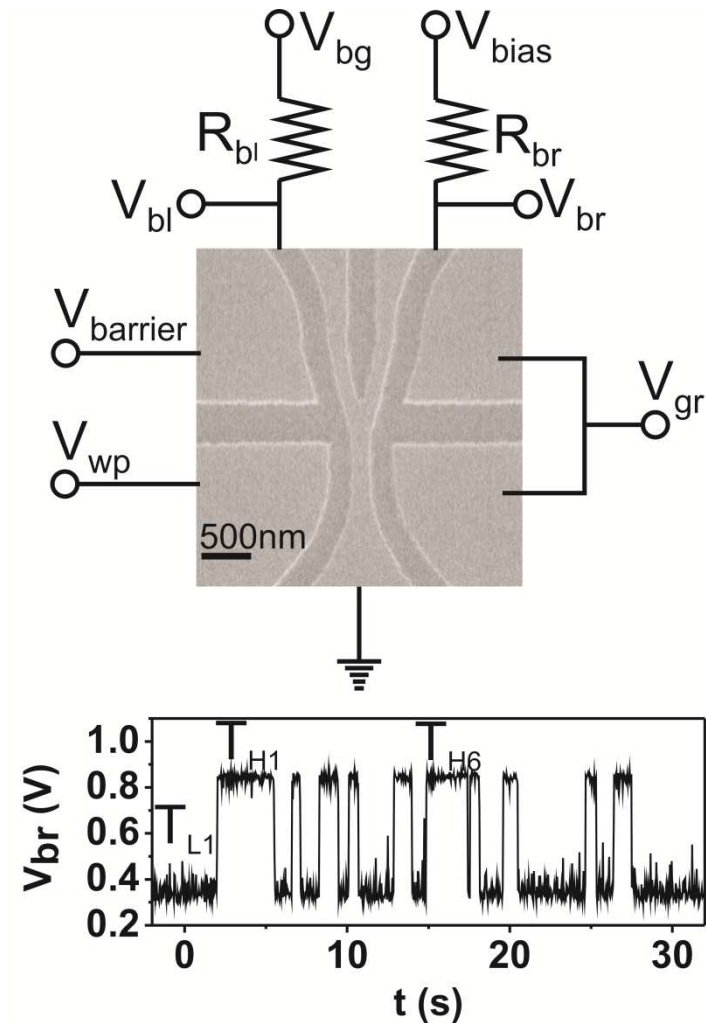


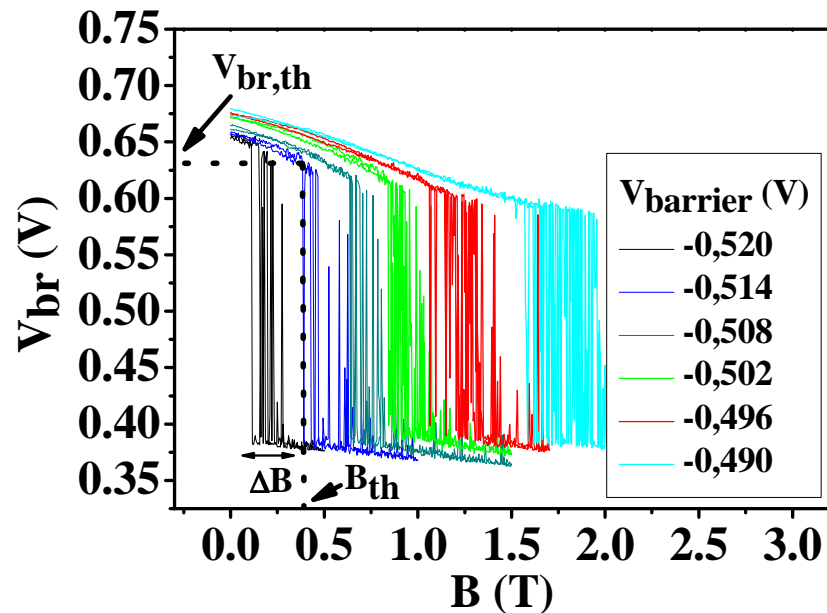
From: L. Gammaitoni and A.D. Bulsara, „Noise Activated Nonlinear Dynamic Sensors“, PRL **88**, 230601-1 (2002).

- The detector is biased in the strongly noise activated regime.
- Switching between V_H and V_L solely controlled by the internal noise.
- Magnetic field is applied perpendicular to the motion of electrons.
- Measure the time spent in each of the two stable states:

$$T_{H,L} = \frac{1}{n_{H,L}} \sum_{i=1}^{n_{H,L}} T_{H_i,L_i}$$

- Output of the detector is the residence time difference: $\Delta T = T_H - T_L$

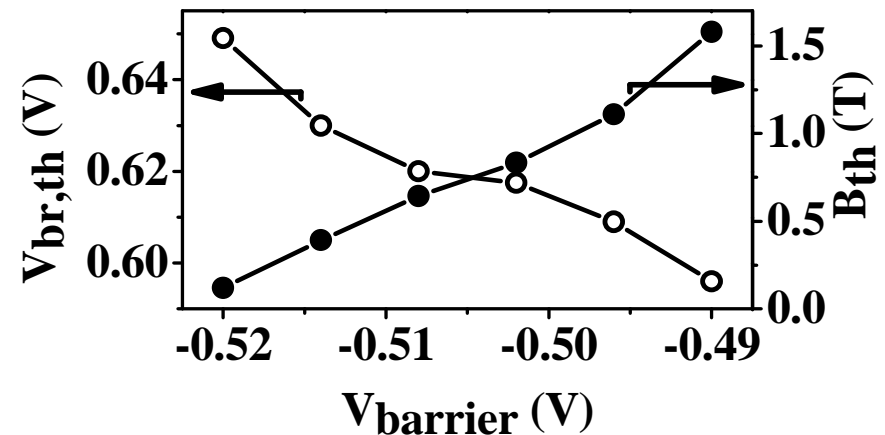




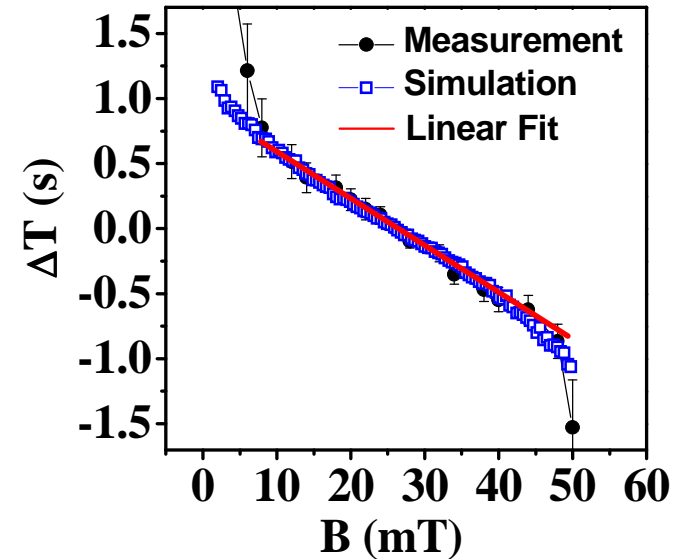
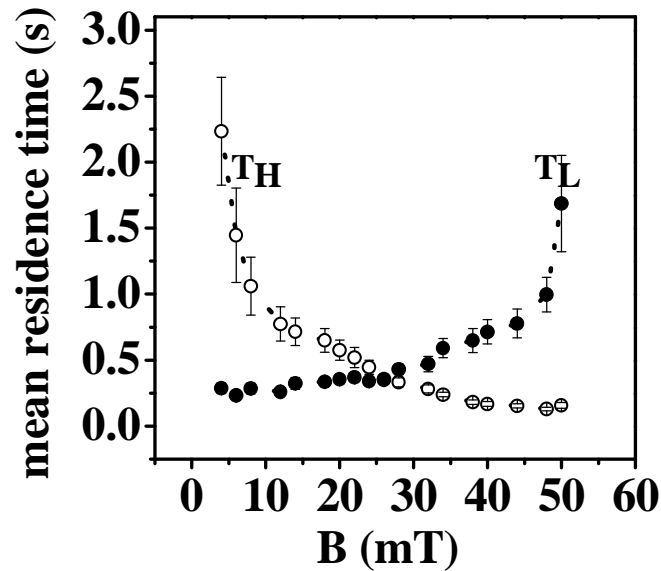
Increasing magnetic field:

- The output V_{br} decreases linearly down to a magnetic field threshold B_{th} .
- Transitions between the two stable states occur within a magnetic field range ΔB .
- The output V_{br} changed its stable state from $V_{br} = V_H$ to $V_{br} = V_L$.

- The magnetic-field induced switching (between V_H and V_L) is associated with an interplay between a scattering asymmetry at the boundaries. [1]



[1] D. Hartmann *et al.*, PHYSICAL REVIEW B **78**, 113306 (2008).



- The residence time T_H (high state) is decreasing and T_L (low state) is increasing with increasing B .
- Output ΔT is a linear function of the magnetic field around the symmetric point $\Delta T = 0$ s.
- Target signal (magnetic field) independent sensitivity.

$$\Delta T(B) = T_0 - cB$$

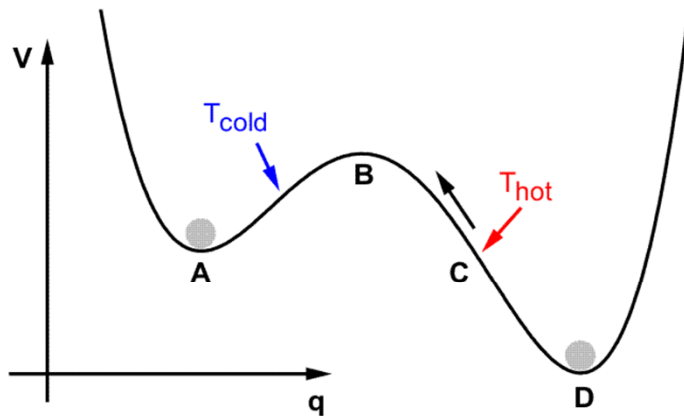
$$S(B) = \frac{\partial \Delta T}{\partial B} = c$$

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Diffusion constants:

$$D = \mu kT_H, \quad q \text{ in } H$$

$$D = \mu kT_L, \quad q \text{ in } L$$



- Double well potential with minima located at A and D.
- D is the energetic favorable point D with $D < A$.
- Consider two temperatures at the slopes T_{hot} and T_{cold} with $T_{\text{hot}} > T_{\text{cold}}$.

For systems subject to thermal noise, the Boltzmann factor is

$$\exp\left(\frac{-V}{kT}\right)$$

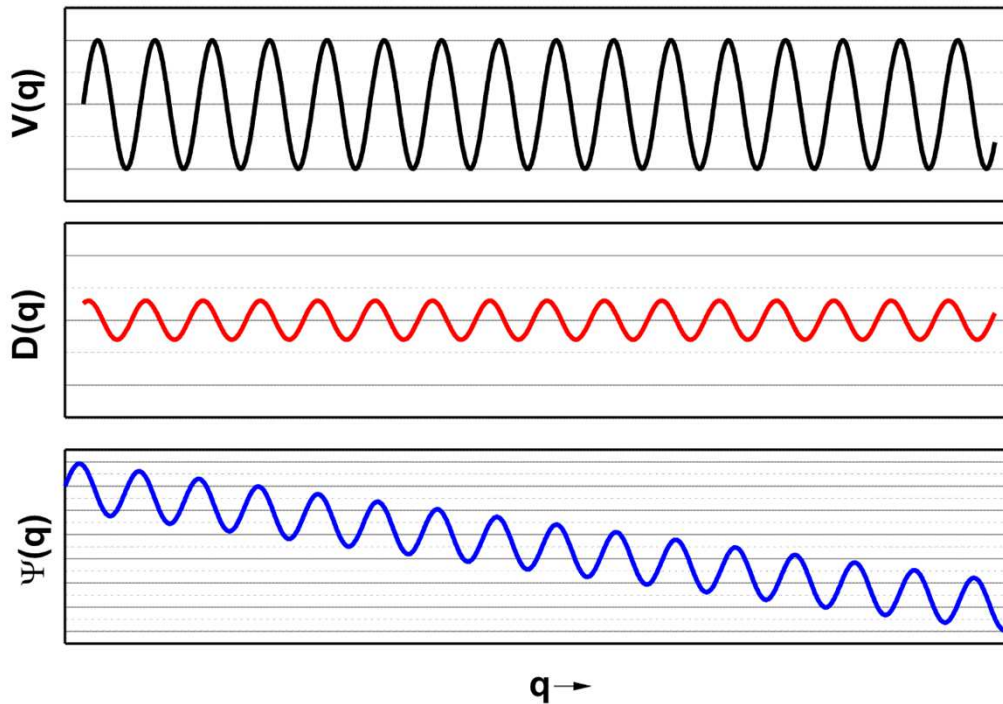
M. Büttiker, Z. Phys. B **68**, 161 (1987).

R. Landauer, J. Stat. Phys. **53**, 233 (1988).

For systems with mobility μ subject to drift and state dependent diffusion the Boltzmann factor is

$$\exp(-\Psi(q))$$

$$\text{with } \Psi(q) = - \int_0^q dp \frac{v(p)}{D(p)}$$



$$V(q) = V(q + 2\pi)$$

$$V(q) = V_0(1 - \cos(q))$$

$$D(q) = D(q + 2\pi)$$

$$D^{-1}(q) = D_0^{-1} (1 - \alpha \cos(q - \phi))$$

$$D_0 = \mu kT$$

$$\Psi(q) = - \int_0^q dp \frac{v(p)}{D(p)}$$

$$\Psi(q) = \Psi(q + 2\pi) + 2\pi \Delta$$

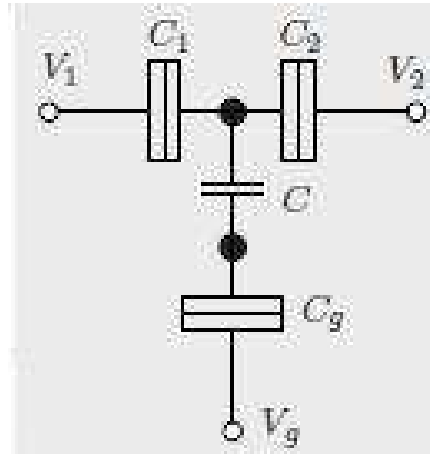
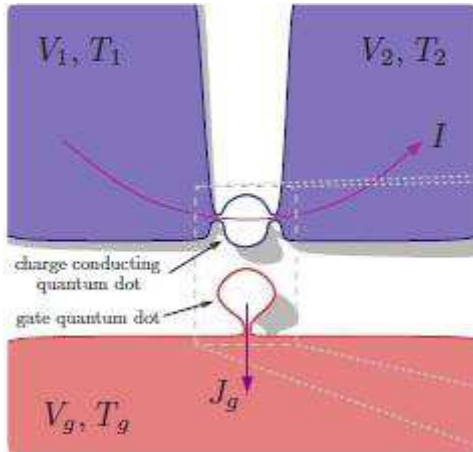
With: $\Delta = \frac{\mu V_0 \alpha}{D_0} \frac{\sin(\phi)}{2}$

M. Büttiker, Z. Phys. B **68**, 161 (1987).

Ya. M. Blanter and M. Büttiker, Phys. Rev. Lett. **81**,
4040-4044 (1998).

$$I_{ov} = \frac{\pi^2 E_0^2 T_1}{\mathcal{L}^2 T_0^2} \exp\left(-\frac{E_0}{T_0}\right) \sin(\varphi)$$

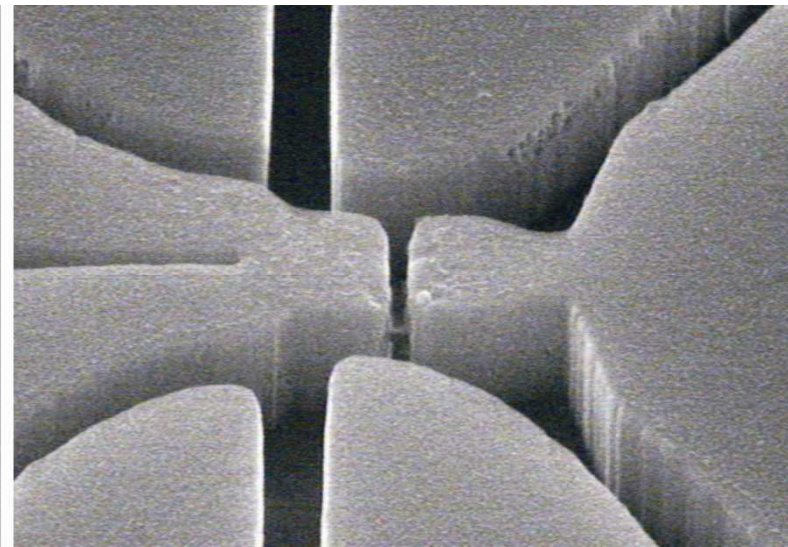
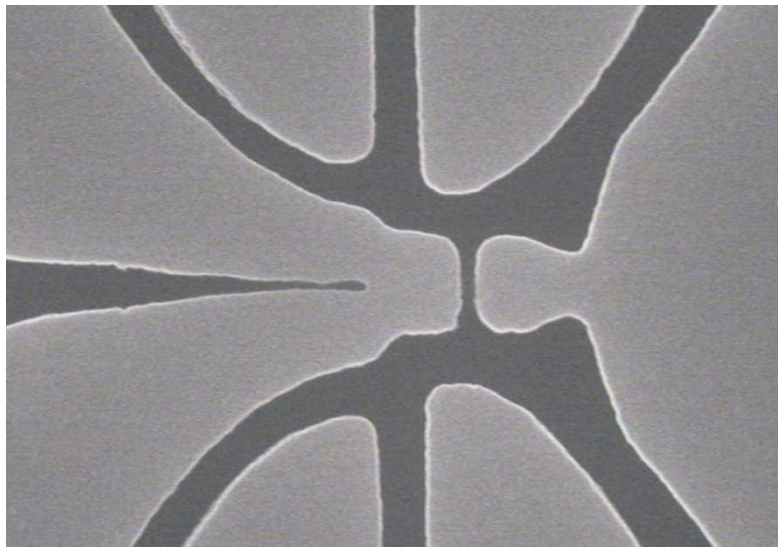
$$I = \frac{\mathcal{Y} T_1}{2m T_0} \exp\left(-\frac{E_0}{T_0}\right) \sin(\varphi)$$



Proposed by R. Sanchez and M. Büttiker, PRB **83**, 085428 (2011).

$$\frac{I}{q} = \frac{J_g}{E_C}$$

Every energy quantum of heat flow gets converted into a quantum of charge flow



- **Growth, fabrication and transport properties of nanoelectronic devices**
 - Samples are based on GaAs/AlAs and grown by molecular beam epitaxy
 - Dry and wet chemical etching is used to define the structures
 - **Universal logic gate switching in resonant tunneling diodes (RTDs)**
 - Two universal logic gates => NOR to NAND for $\Delta V_{ac} \sim 0.1 \text{ mV}$
 - Logic stochastic resonance (LSR) with $P_{noise} \sim \text{nW}$
 - **Stochastic resonance in nanoelectronic devices**
 - SR @ $f = 1 \text{ Hz}$ in electron waveguides: Tuning the periodic forcing
 - SR @ $f = 500 \text{ Hz}$ in RTDs for ac and periodic optical modulation: $P_{light} = 160 \text{ nW}$
 - **Noise activated nonlinear dynamic sensors**
 - Magnetic field sensor based on (bistable) electron waveguides
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- www.nanopwr.eu

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Many thanks for your attention!
